

Neutron cross section measurements in the resonance region at GELINA



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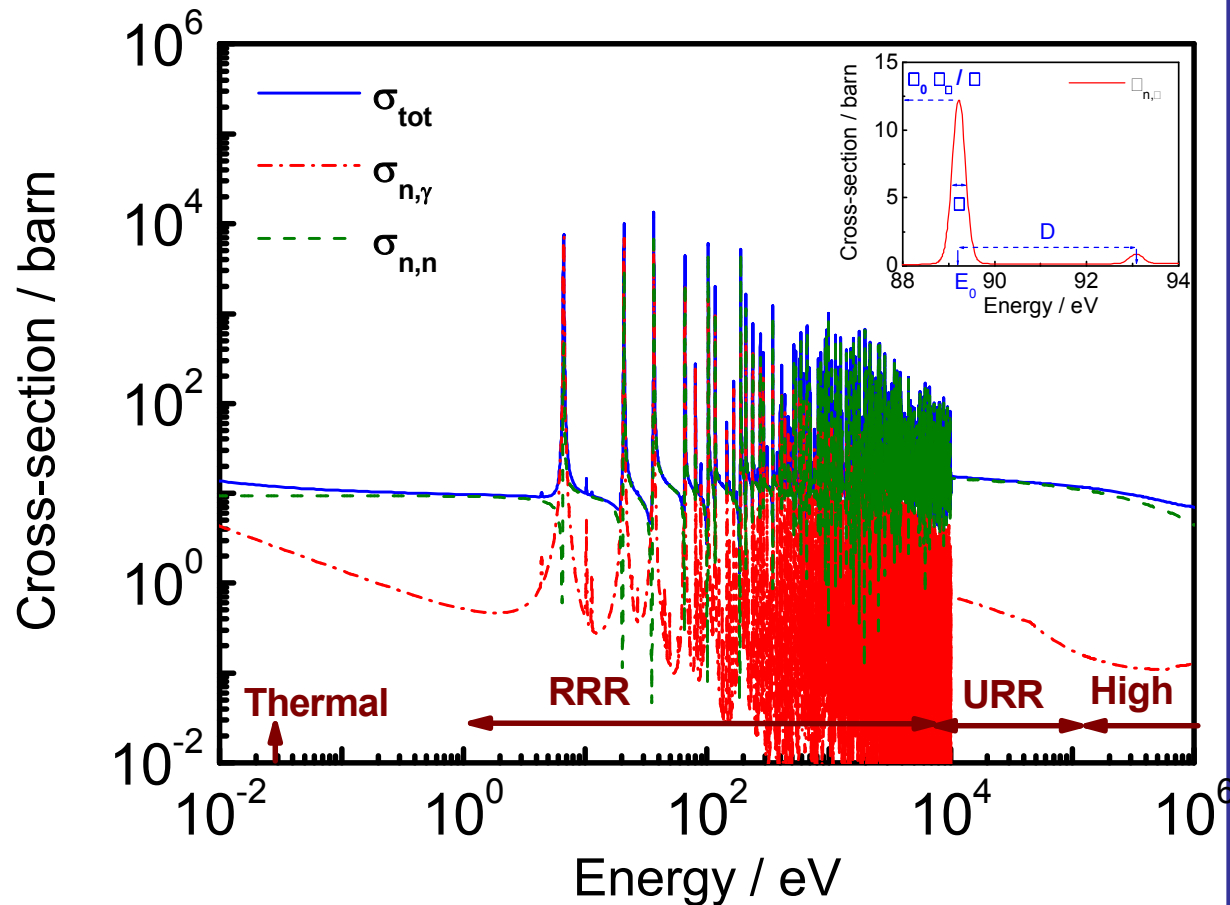
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Measurements**

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<http://irmm.jrc.ec.europa.eu/>

<http://www.jrc.ec.europa.eu/>

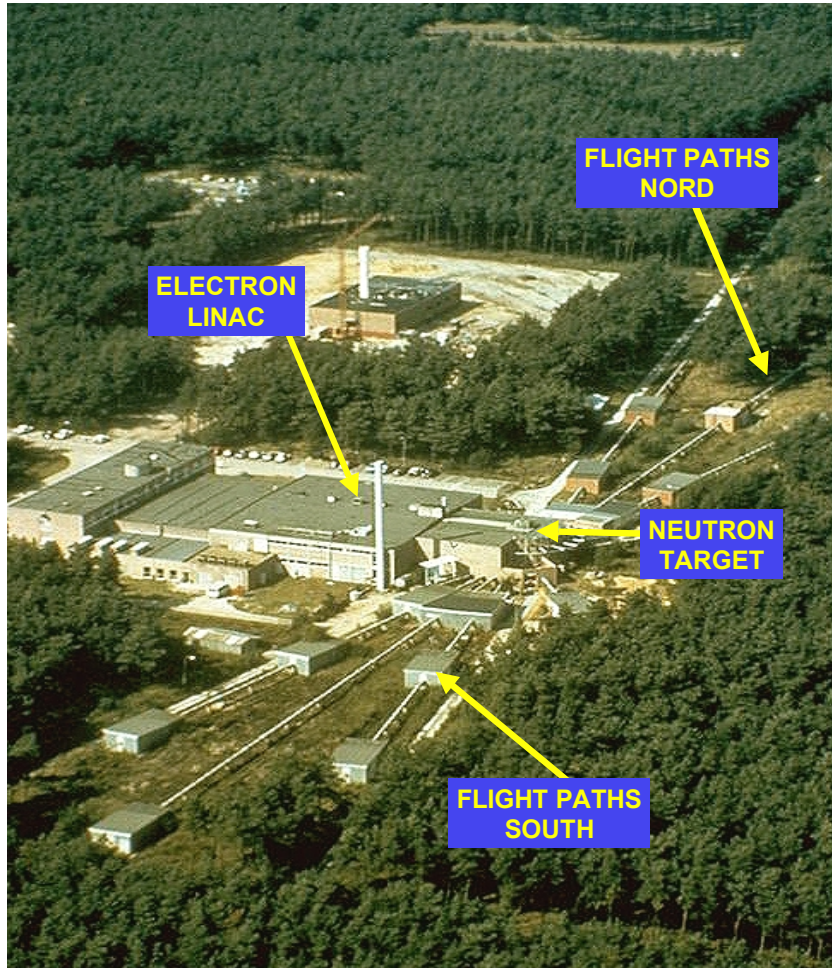
- **Introduction**
- **GELINA – TOF Facility**
- **Cross Section Measurements**
- **Data analysis**
- **Resonance analysis**
- **Results: Cd, Rh, Au**



$${}^{238}\text{U}(n,\text{tot}) = {}^{238}\text{U}(n,n) + {}^{238}\text{U}(n,\gamma)$$

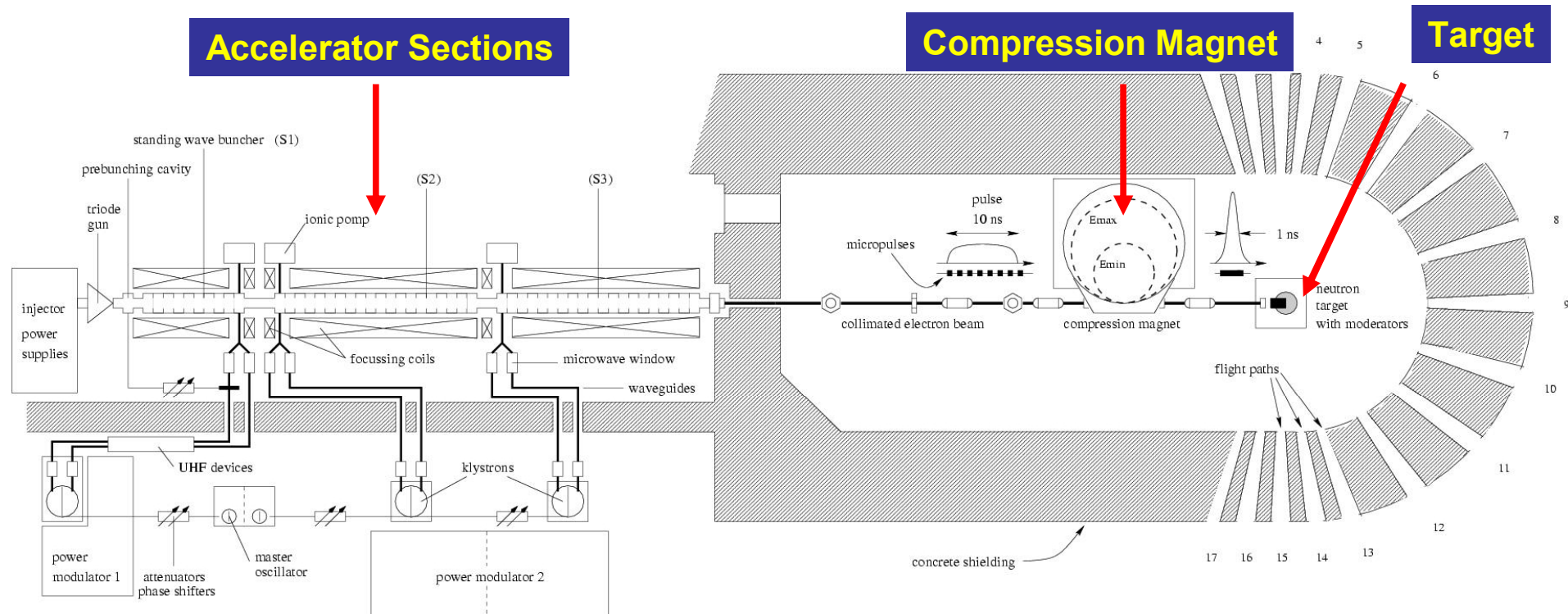
$$\sigma_{\text{tot}} = \sigma_n + \sigma_\gamma$$

- **Thermal**
- **Resonance Region : $D > \Gamma$**
 - Resolved Resonance Region : $\Delta_R < D$
 - Unresolved Resonance Region : $\Delta_R > D$
- **High Energy Region : $D < \Gamma$**



- Time-of-flight facility
- Pulsed white neutron source
($1 \text{ meV} < E_n < 20 \text{ MeV}$)
- Multi-user facility with 10 flight paths
(10 m - 400 m)
- The measurement stations have special equipment to perform:
 - Total cross section measurements
 - Partial cross section measurements

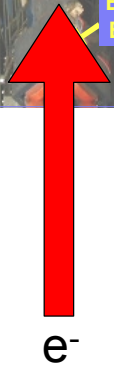
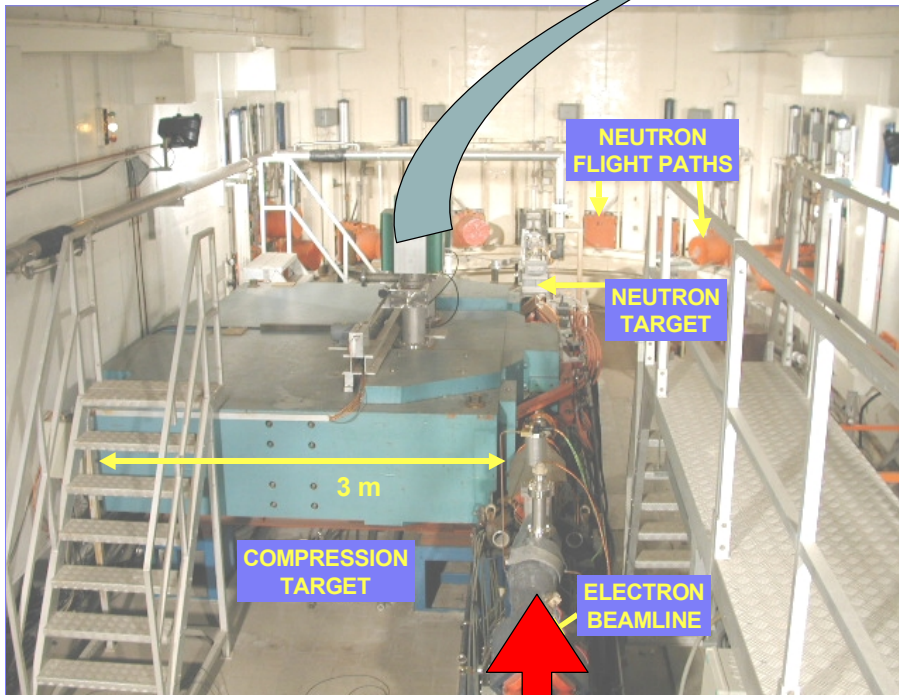
Pulse Width : 1ns
Frequency : 50 – 800 Hz
Average Current : 4.7 – 75 μA
Neutron intensity : $1.6 \cdot 10^{12}$ – $2.5 \cdot 10^{13}$ n/s



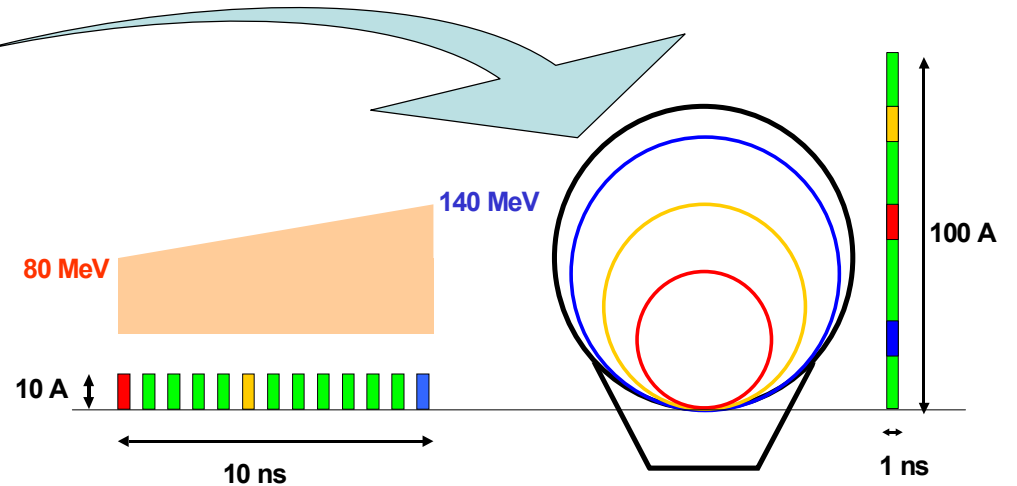
Normal Operation Parameters

Average Current : 75 μA
Average Electron Energy : 100 MeV
Mean Power : 7.5 kW

Frequency : 800 Hz
Pulse Width : 1 ns
Neutron Intensity : 2.5×10^{13} n/s



e^-



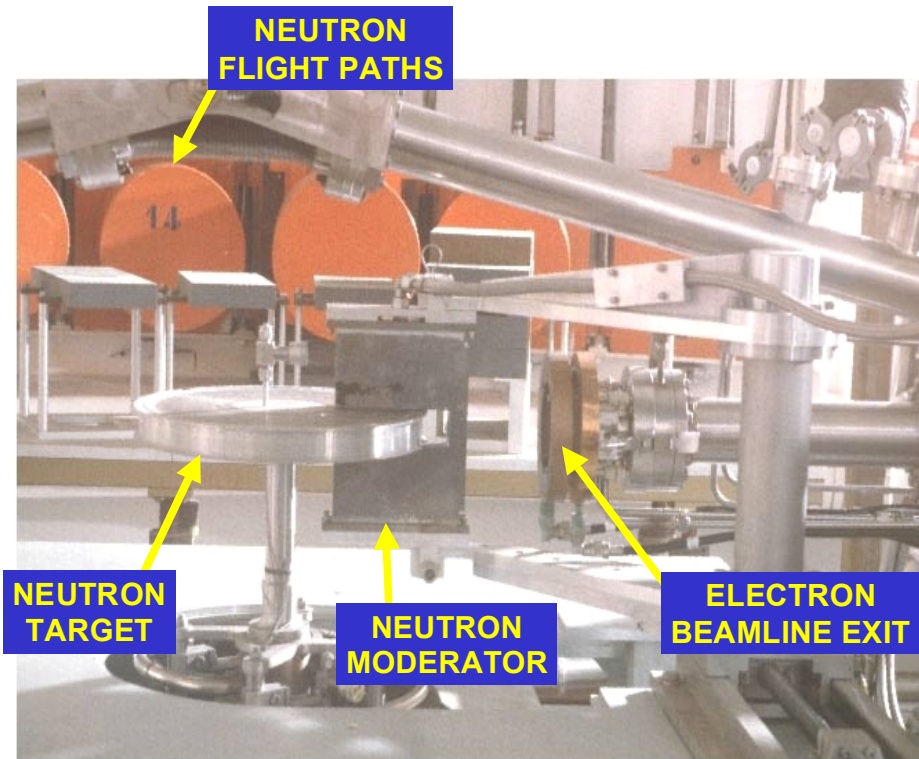
$$B\rho = \frac{p}{q}; E \cong pc; q=e$$

$$\Rightarrow \rho = \frac{1}{B} \frac{E}{qc}$$

$$\Rightarrow B = \frac{2\pi}{qc^2} \frac{\Delta E}{\Delta\tau}$$

$$\begin{aligned} \Delta E &= 60 \text{ MeV} \\ \Delta\tau &= 10 \text{ ns} \end{aligned}$$

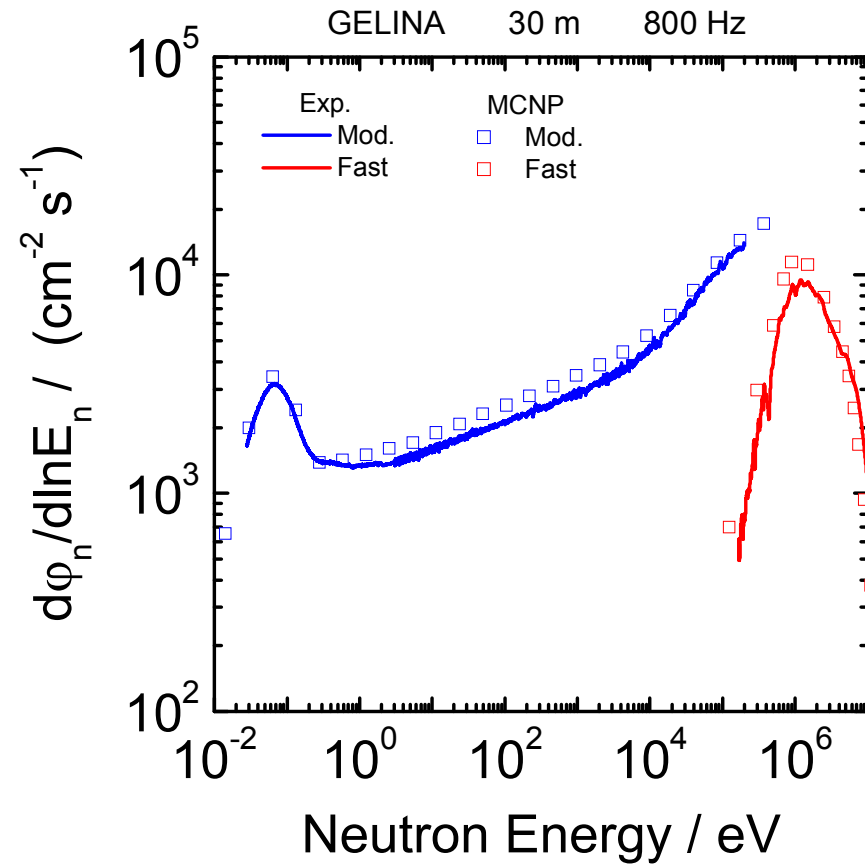
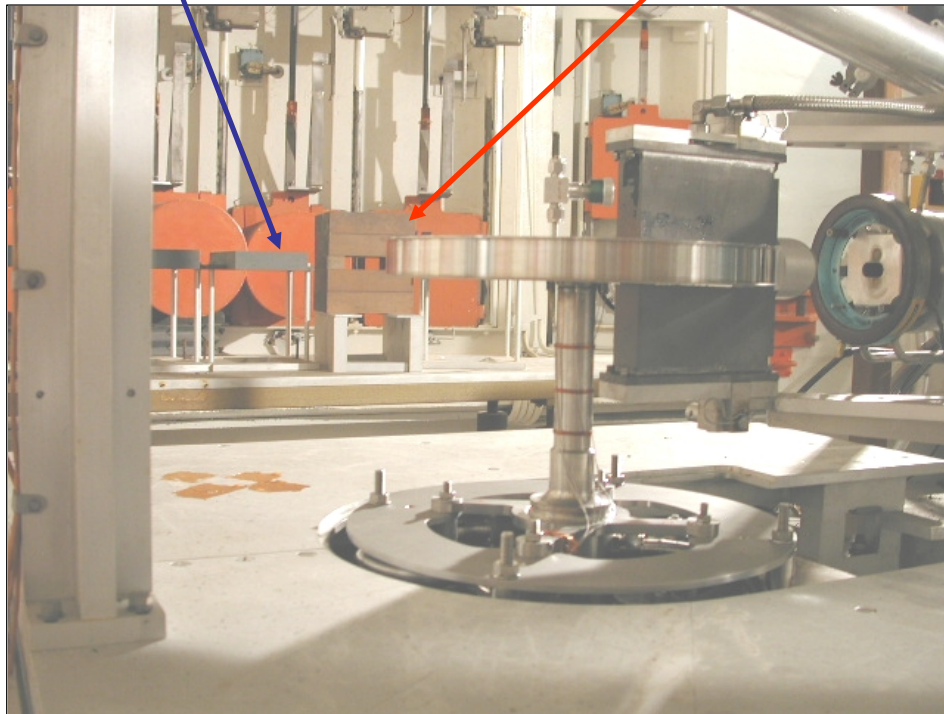
→ compressed pulse length $\sim 1 \text{ ns}$

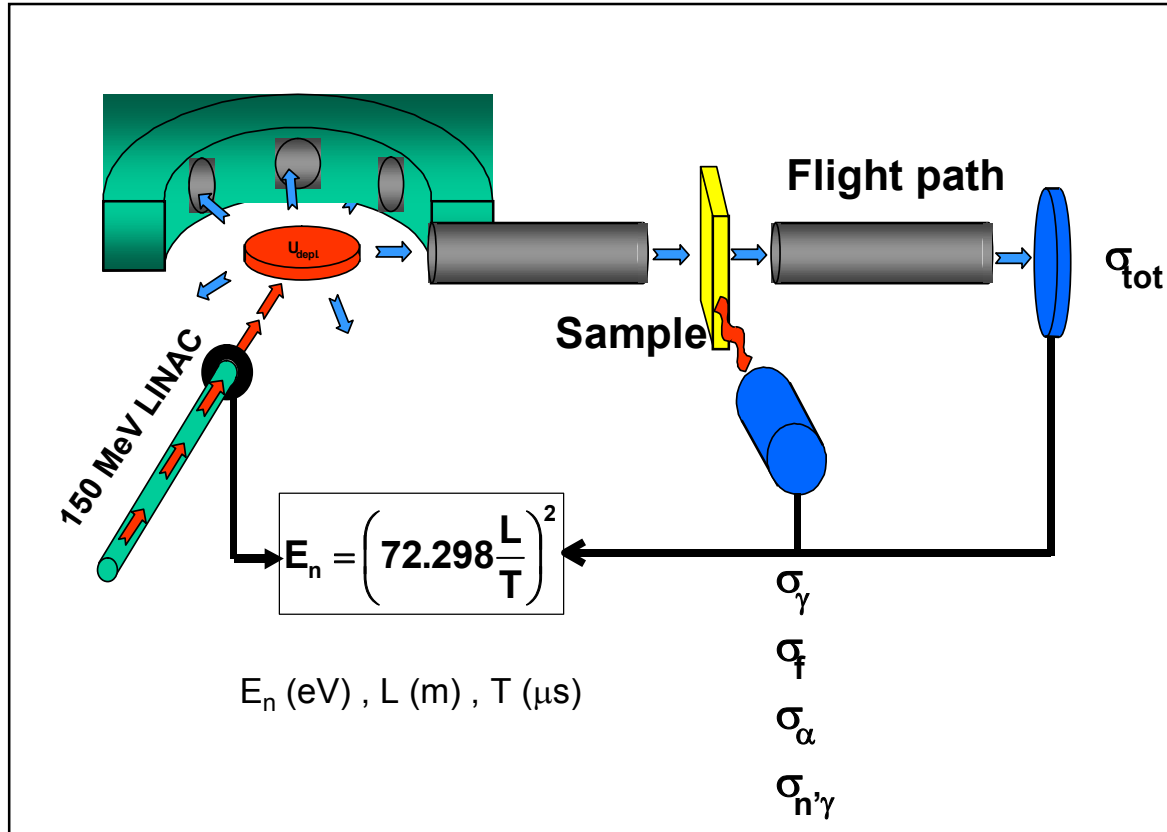


- e^- accelerated to $E_{e^-,max} \approx 140 \text{ MeV}$
- (e^-, γ) Bremsstrahlung in U-target (rotating & cooled with liquied Hg)
- (γ, n) , (γ, f) in U-target
- Low energy neutrons by water moderator in Be-canning

SHIELDING for
MODERATED SPECTRUM

SHIELDING for
FAST SPECTRUM





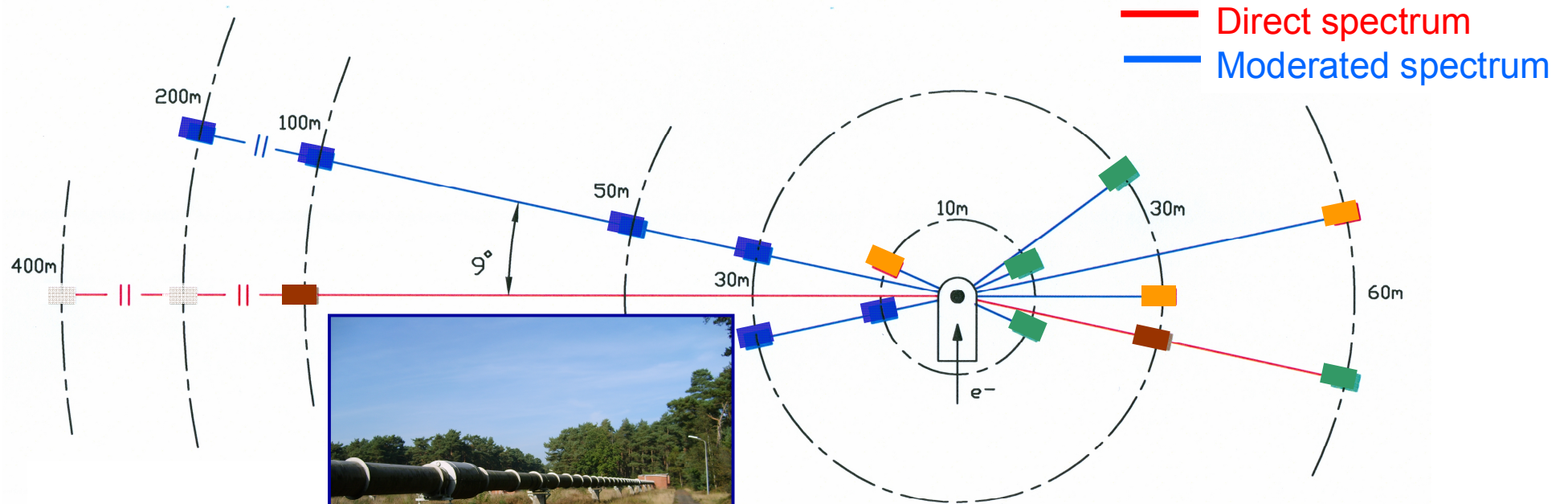
- Neutron Flux $\Rightarrow L \searrow$

$$\phi_n(L) \propto \frac{1}{L^2}$$

- Resolution $\Rightarrow L \nearrow$

$$\frac{\Delta E_n}{E_n} = \frac{1}{L} \sqrt{\frac{E_n \Delta T^2}{\alpha} + \Delta L^2}$$

Optimization of flight path length in view of neutron flux and resolution !



- (n,γ)
- (n,tot)
- (n,f) and (n,cp)
- (n,n'γ)

- NIM A, 577, 626 (2007)
- NP A 773, 173 (2006)
- NSE 156, 211 (2007)
- NP A 786, 1 (2007)

- Ge-detectors (L = 10 m)
- C₆D₆ detectors (L = 10, 30 and 60 m)

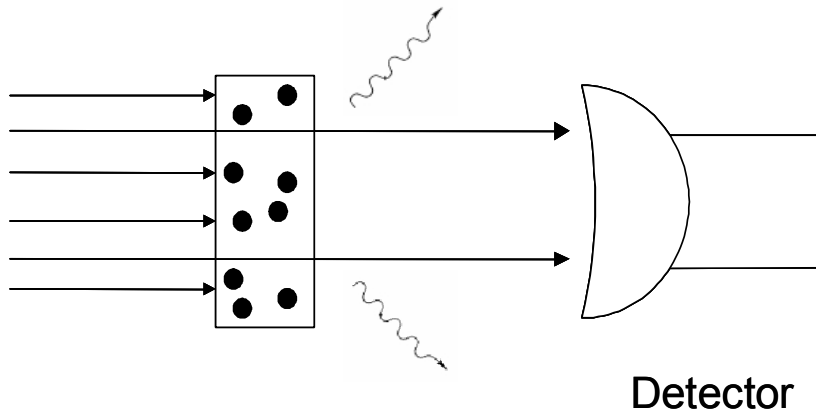
- ⁶Li glass -detectors (L = 25 and 50m)

Transmission : σ_{tot}

$$T = \frac{C_{in}}{C_{out}} \cong e^{-n\sigma_{tot}}$$

C_{in} = sample in

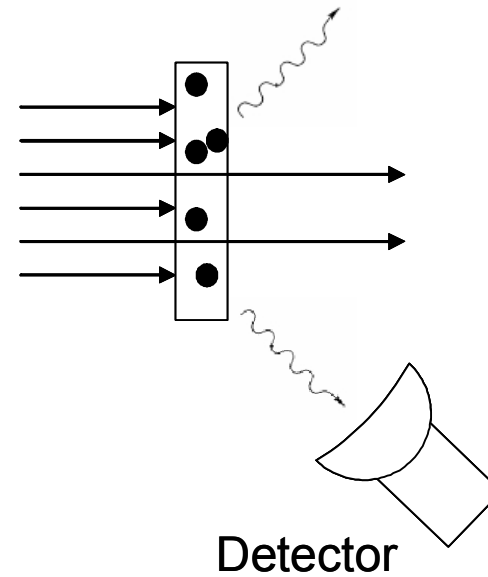
C_{out} = sample out



Partial cross section: $\sigma_{(n,\gamma)}$, $\sigma_{(n,f)}$, $\sigma_{(n,n')}$

$$C_r \propto \varepsilon_r n \sigma_r \phi$$

(thin sample)



Transmission

$$T = \frac{C_{\text{in}}}{C_{\text{out}}} \cong e^{-n\sigma_{\text{tot}}}$$

- Incoming flux cancels
- Detection efficiency cancels
- Good geometry

⇒ Direct relation between T and σ_{tot}

Partial

$$C_r = \varepsilon_r Y_r A_r \phi_r$$

- ϕ_r Neutron Fluence Rate
- ε_r Detection Efficiency
- A_r Effective area
- Y_r Reaction Yield

(Beam Fraction Undergoing the partial Reaction)

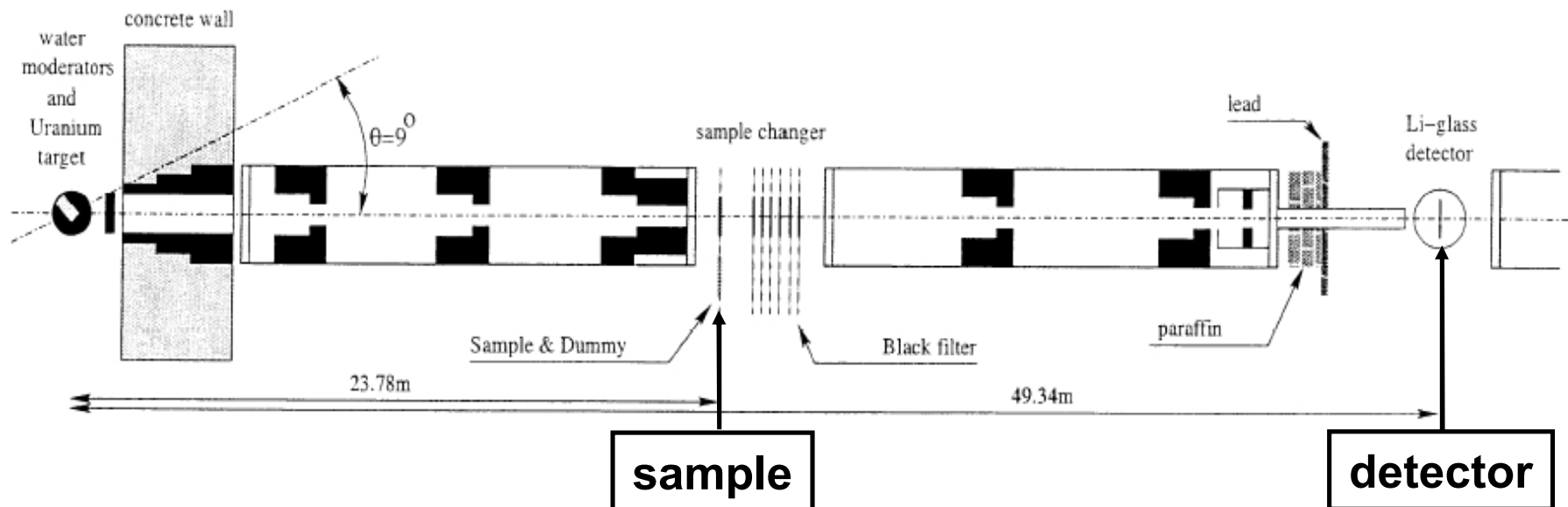
⇒ Complex relation between C_r and Y_r
 Y_r related to σ_r and σ_{tot}

$$T \cong e^{-n\sigma_t}$$

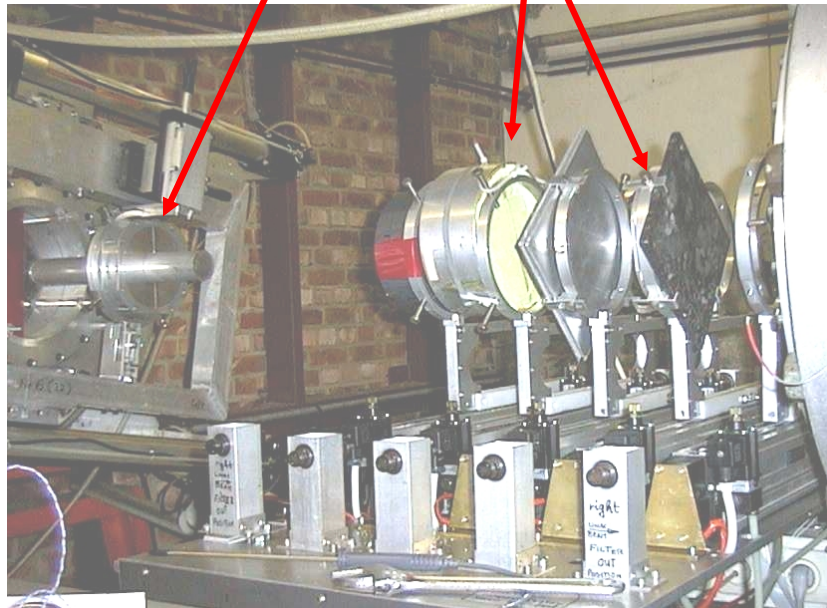
Good Geometry

- All detected neutrons have traversed the sample
- Neutrons scattered in the target do not reach the detector

Importance of collimation



Sample & Background Filters

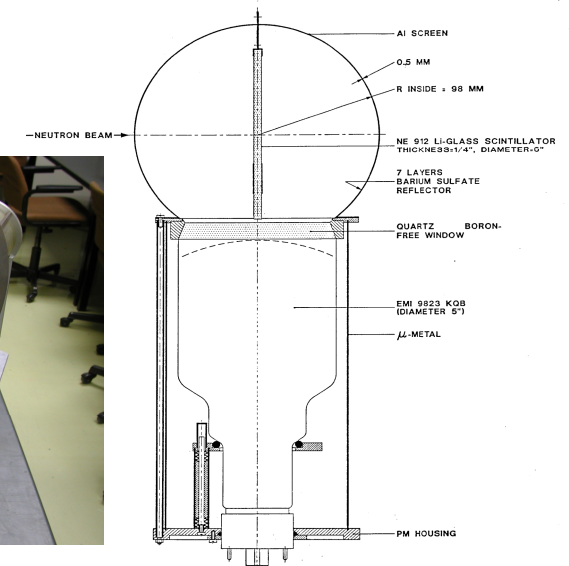
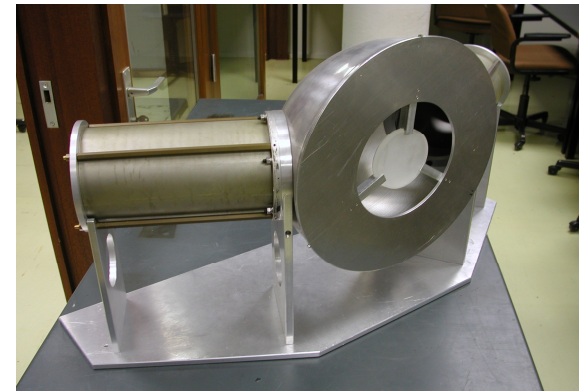


Detector stations

Moderated : L = 30 m, 50 m, (100 m, 200 m)

Fast : L = 400 m

Detector

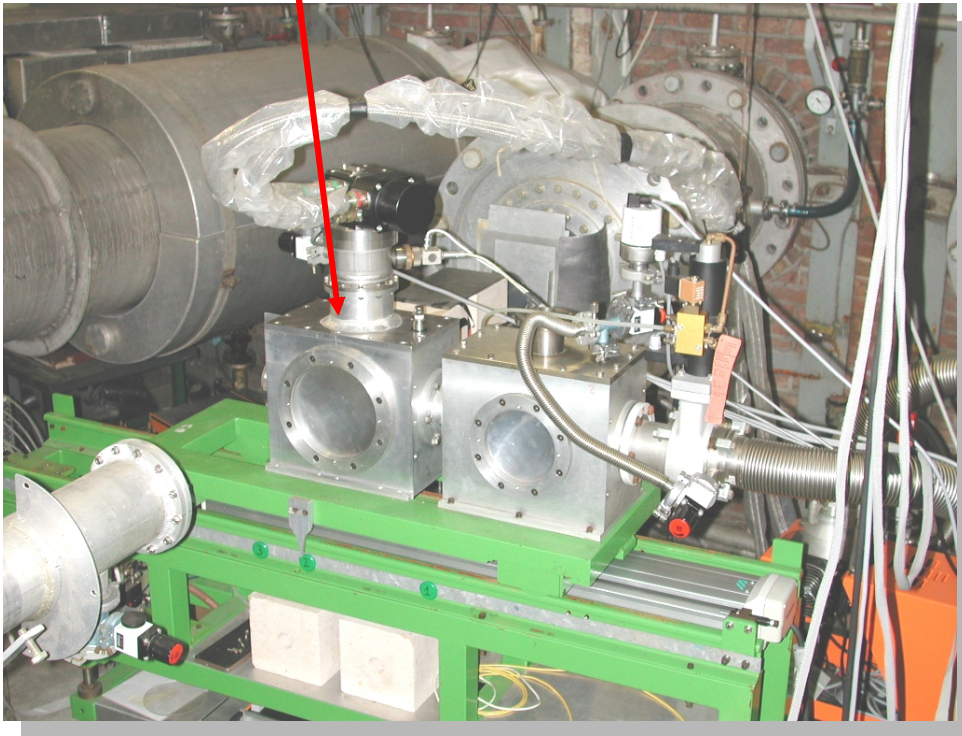


BRUSARZ

Low energy : ${}^6\text{Li}(n,t)\alpha$ Li-glass

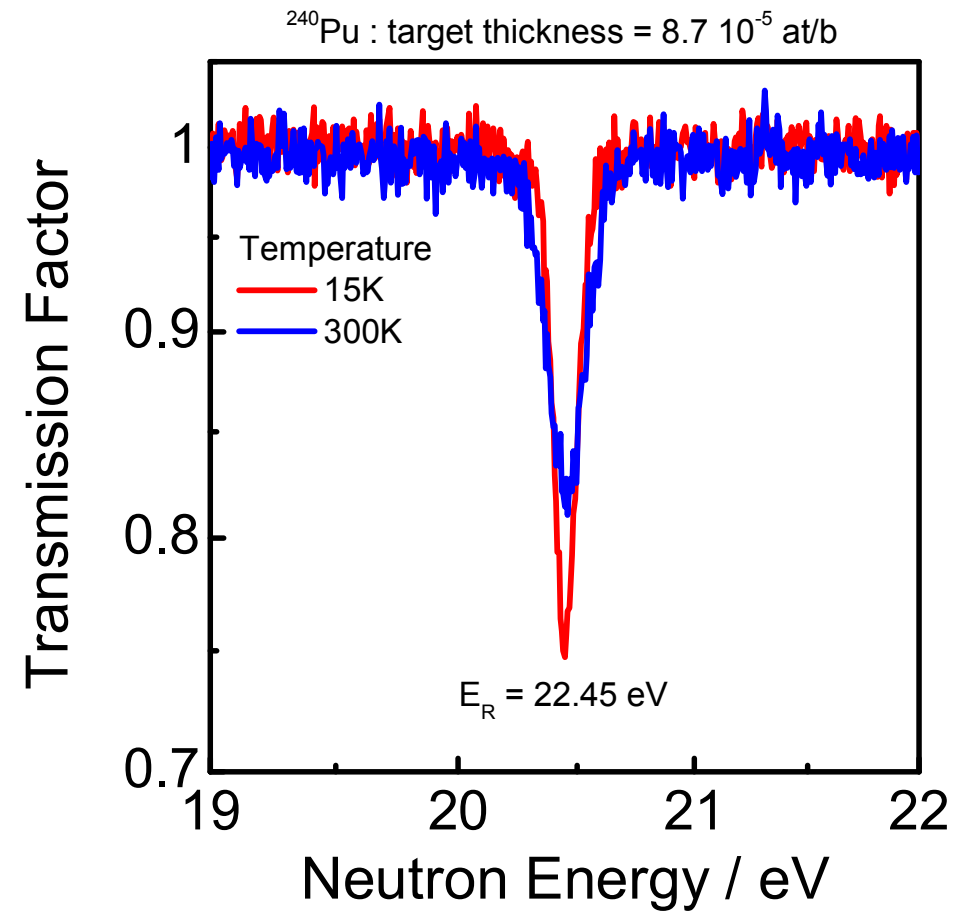
High energy : $\text{H}(n,n)\text{H}$ Plastic scintillator

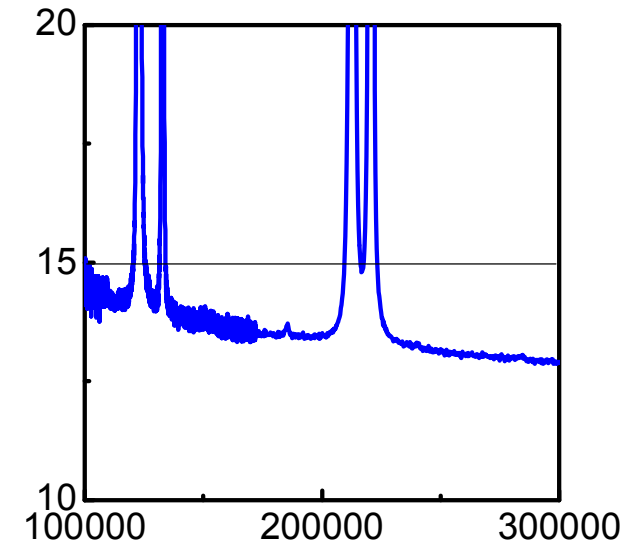
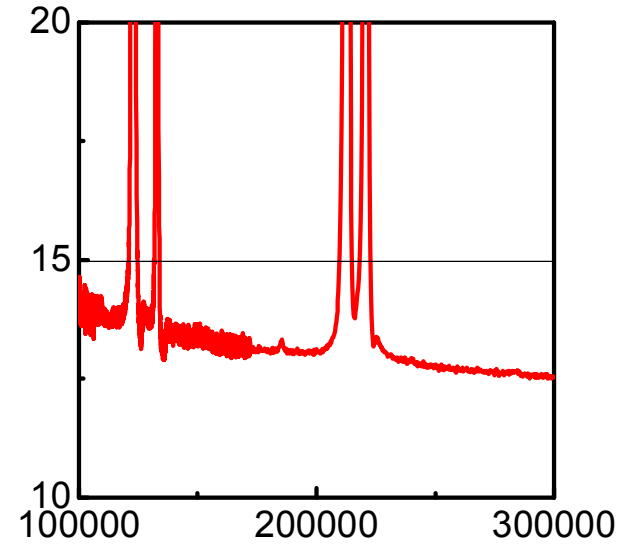
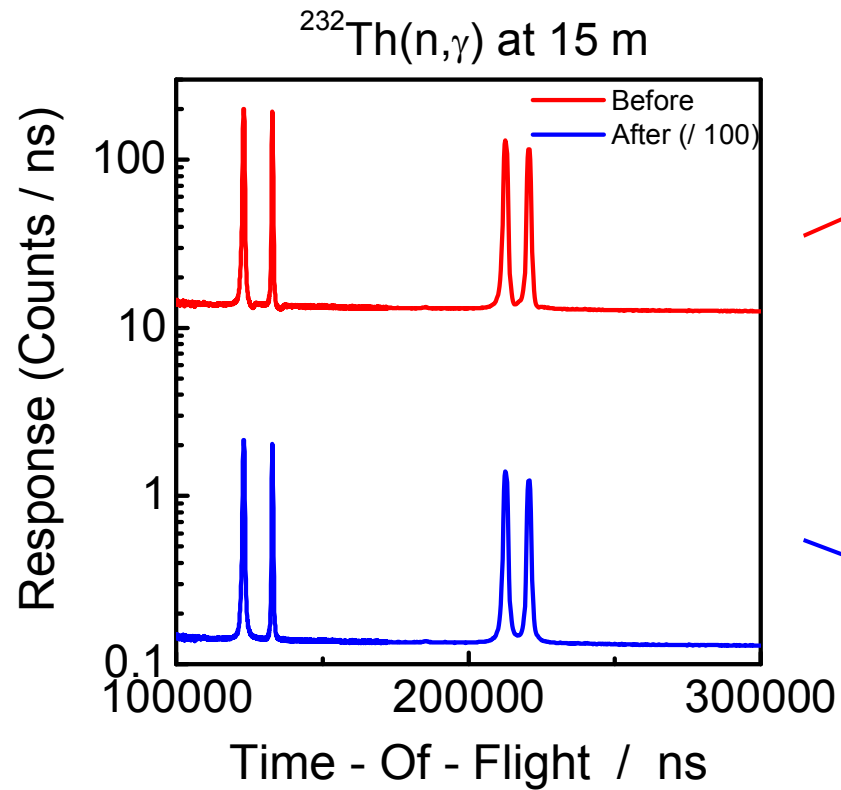
Cryostatic sample changer

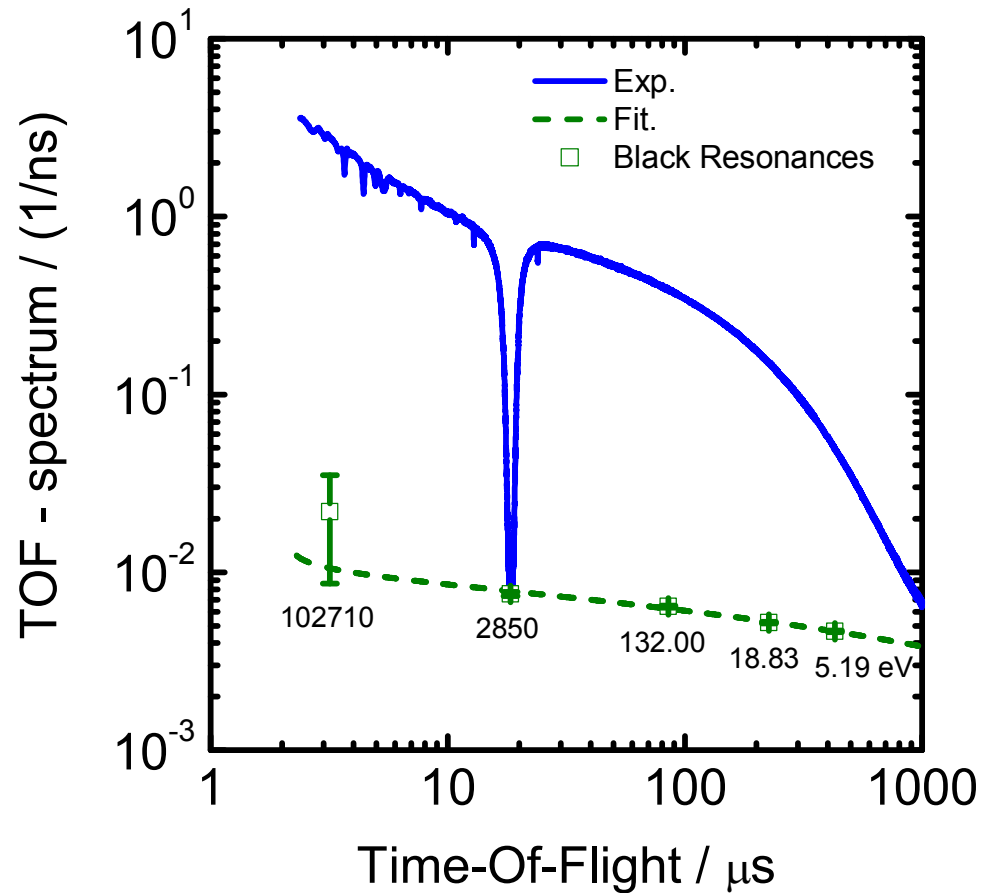


Detector station at : L = 30m

Temperatures from 10 K - 300 K



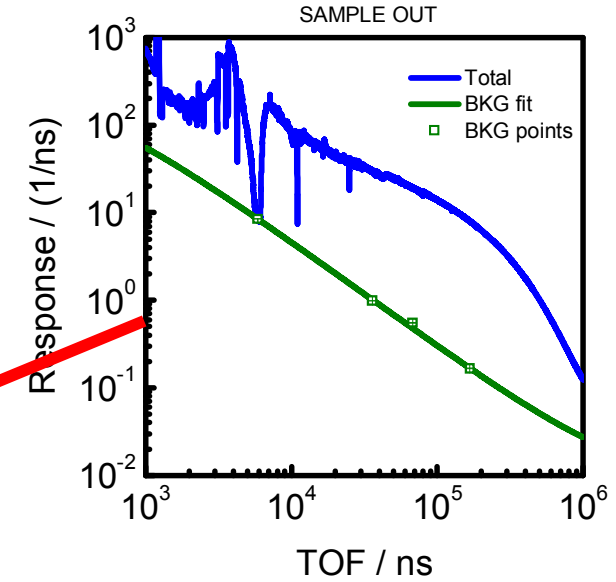
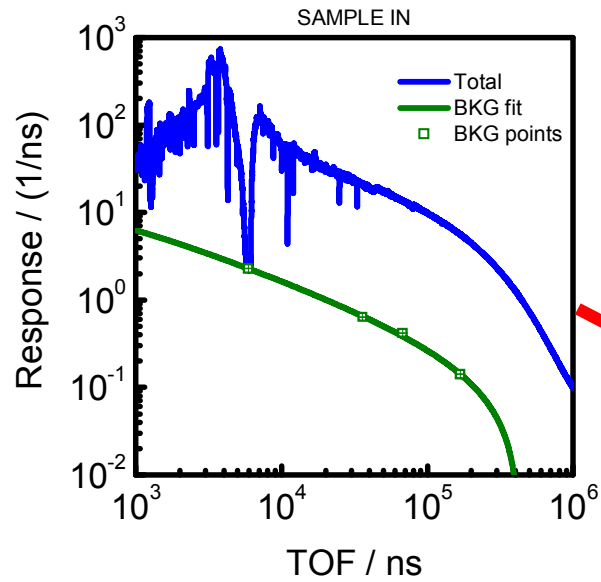




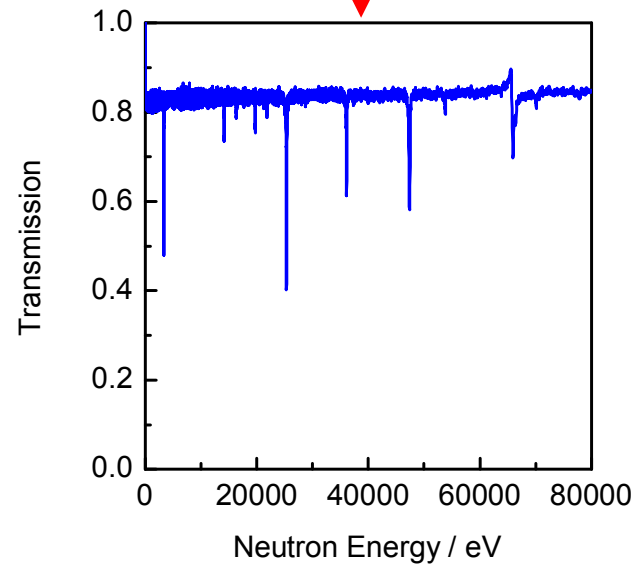
Black Resonance Filters

Element	Energy / eV
Ag	5.19
W	18.83
Co	132.00
Na	2850.00
S	102710.00

A constant and low background during the whole measurement campaign is mandatory !

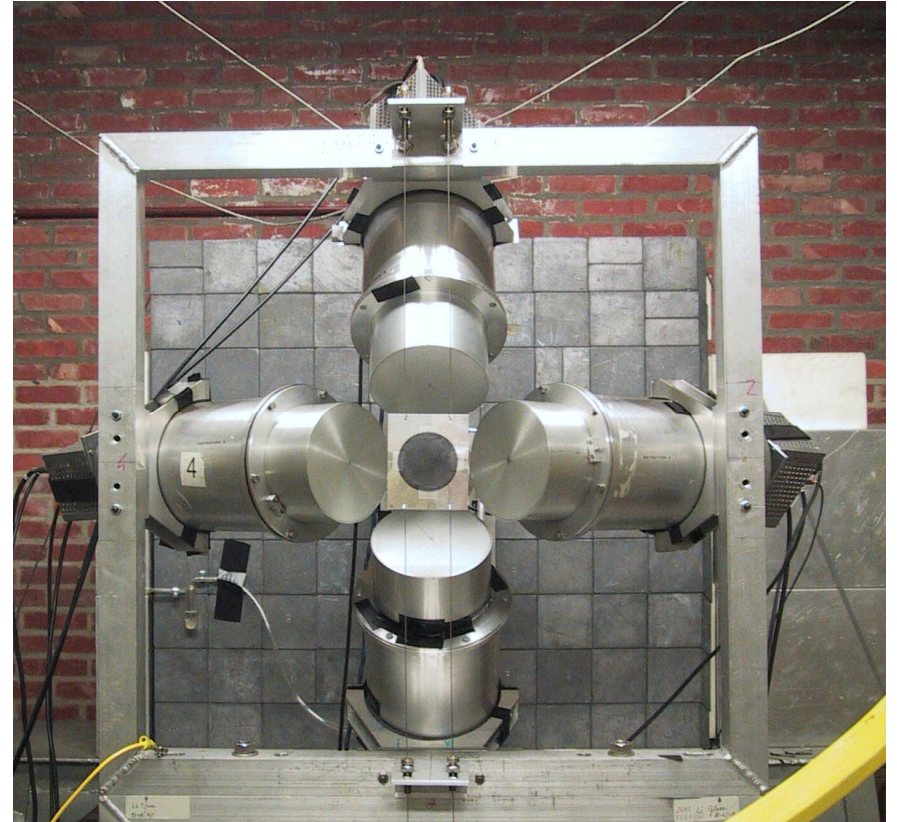


$$T = \frac{C'_{in} - B'_{in}}{C'_{out} - B'_{out}}$$



Total energy detection

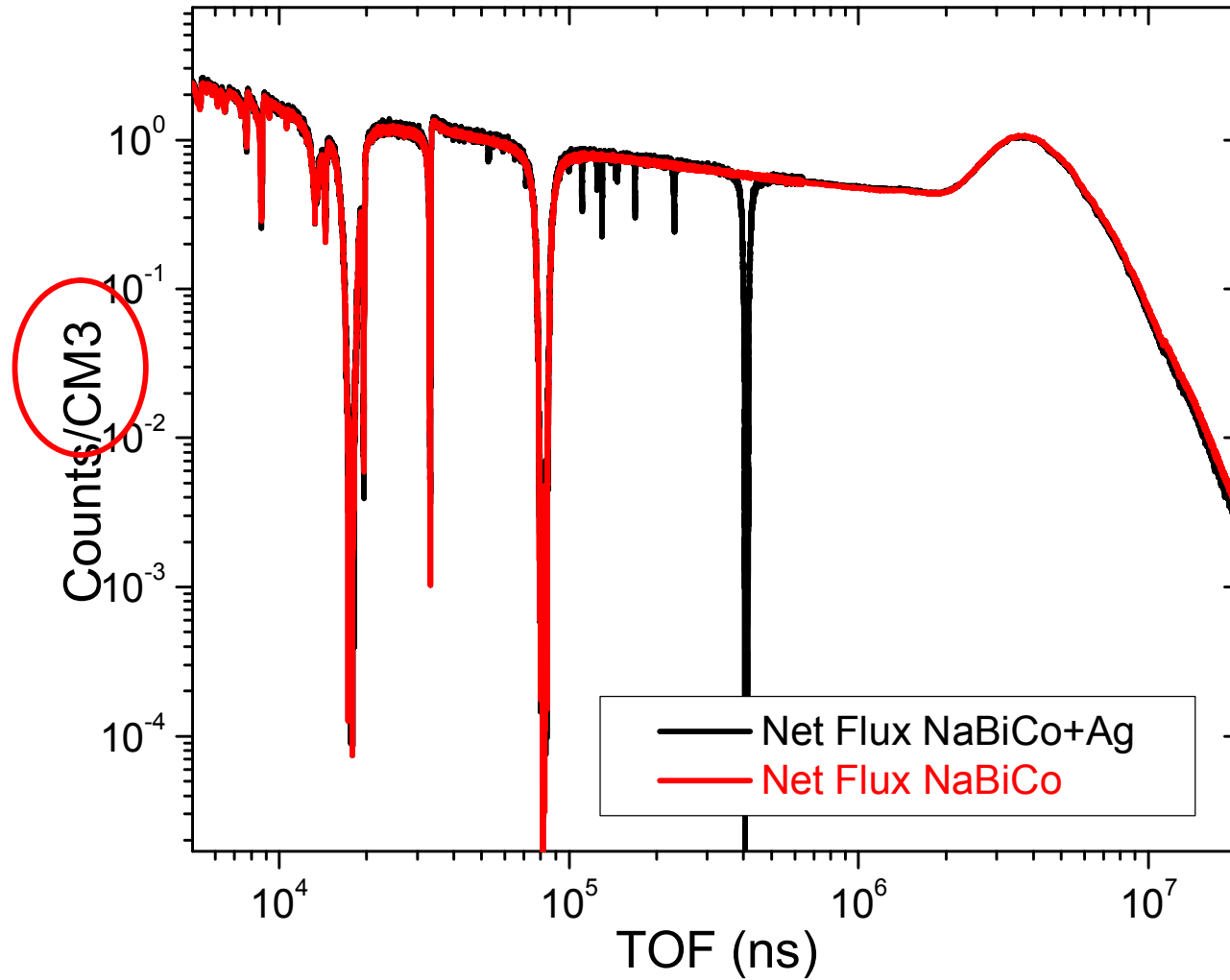
- **C₆D₆ liquid scintillators**
 - 125°
 - PHWT $\int R(E_d, E_\gamma) WF(E_d) dE_d = kE_\gamma$
- **Flux measurements (IC)**
 - ¹⁰B(n,α)
 - ²³⁵U(n,f)

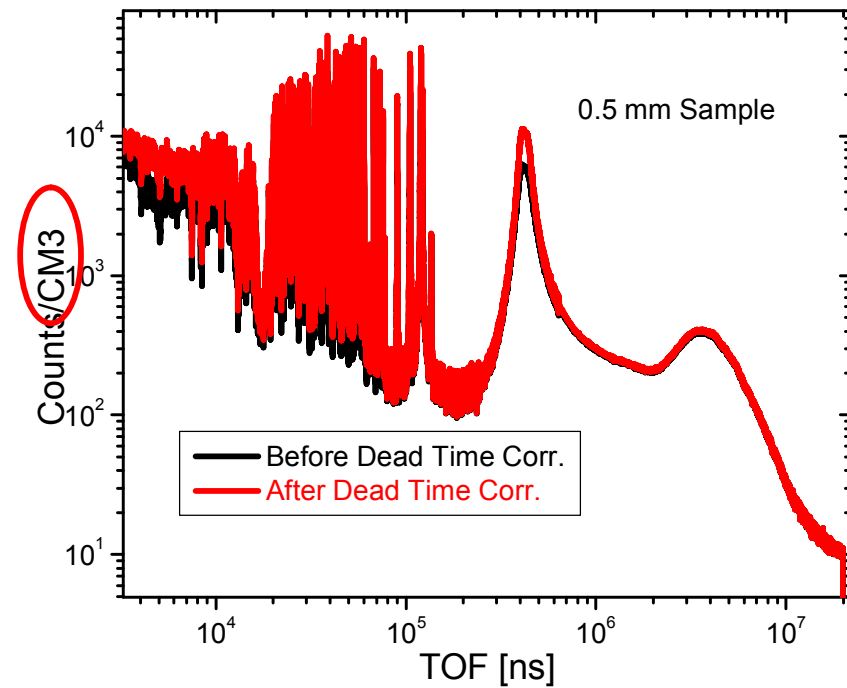
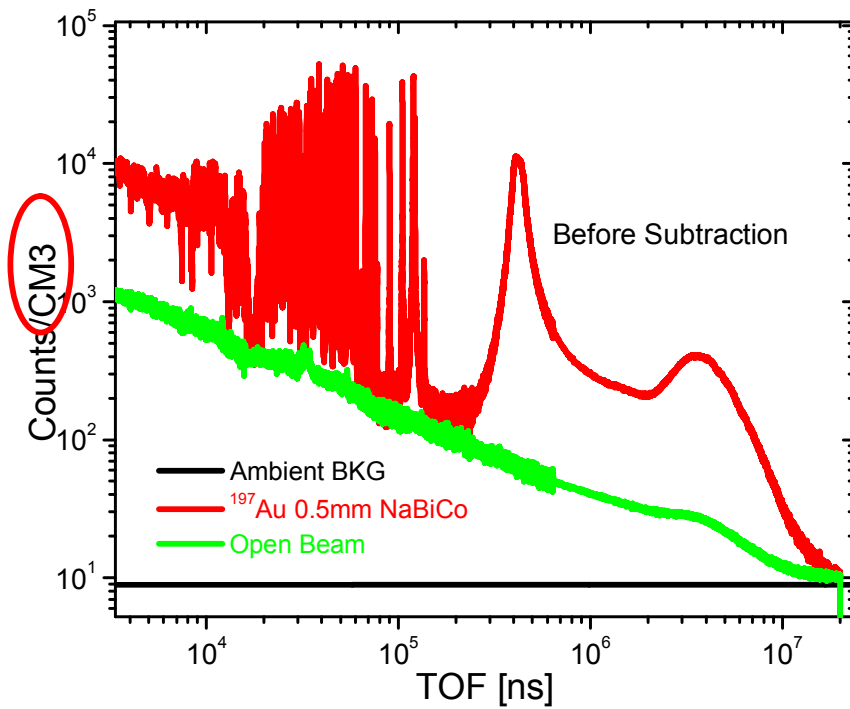


$$Y_{n,\gamma} = \sigma_\varphi \frac{C_w - B_w}{C_\varphi - B_\varphi}$$

L = 10 m, 30 m and 60 m

Borella et al., NIM A, 577, 626 (2007)
 Borella et al., PR C 76, 014605 (2007)

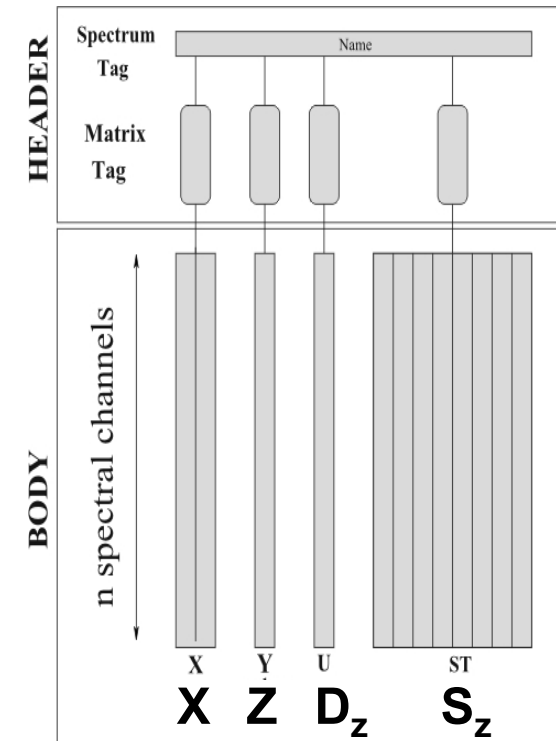




- Transform count rate spectra into observables (transmission factors, partial reaction yields)
- Full propagation of uncertainties starting from counting statistics
- Output: complete covariance matrix
- Due to the special format used in AGS:
 - We reduce space for data storage (EXFOR)
 - We can verify and document the various sources of uncertainties in each step of the reduction process

Observable Z (dimension n) with k sources of correlated uncertainties

$$C_Z = D_Z + S_Z S_Z^T$$



D_Z : uncorrelated part
n values

S_Z : correlated part
dim. (n x k)

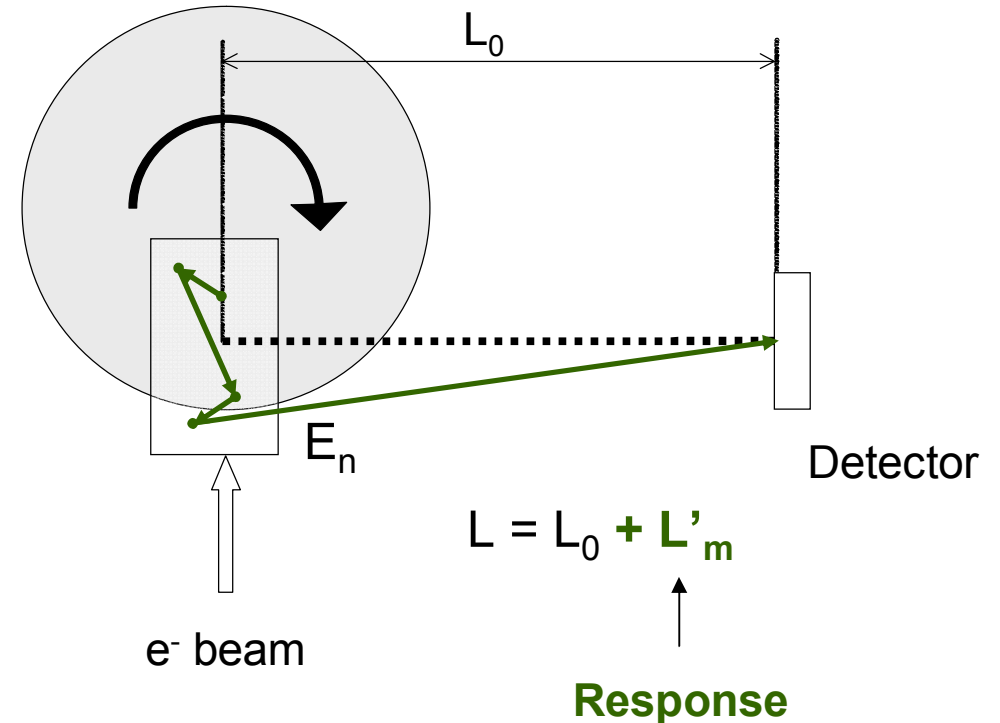
$$R_n = \frac{\Delta E_n}{E_n} = 2 \sqrt{\frac{\Delta T^2}{T^2} + \frac{\Delta L^2}{L^2}}$$

• ΔT

- Initial burst width (~ 1 ns)
- Time jitter of detector ($\sim 1 - 40$ ns)
- Electronics (~ 0.5 ns)

• ΔL

- **Neutron transport in Moderator and U-target**
- Neutron transport in the detector or sample



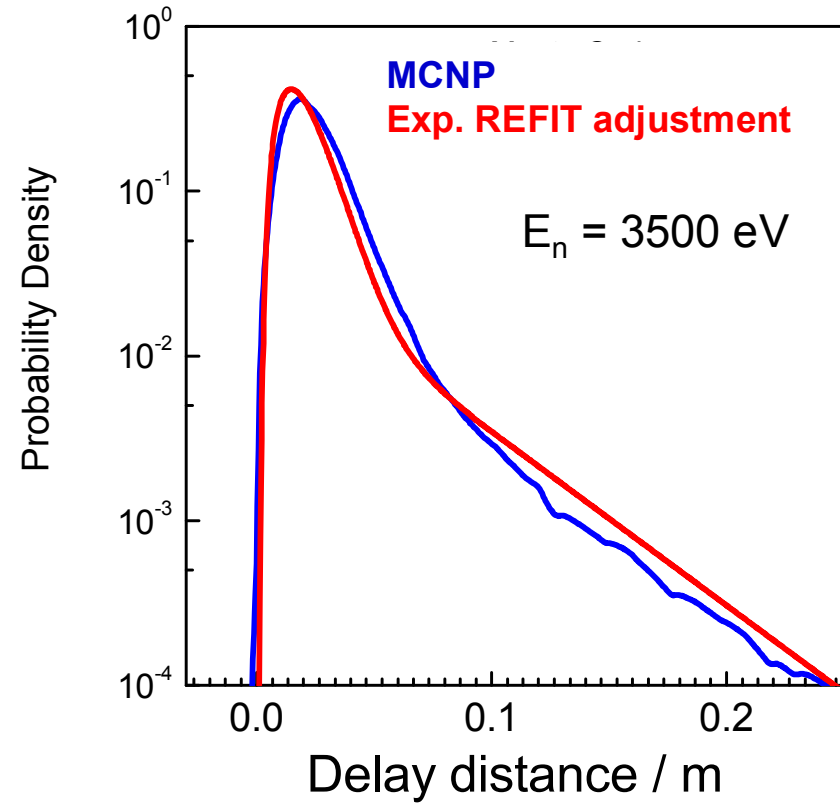
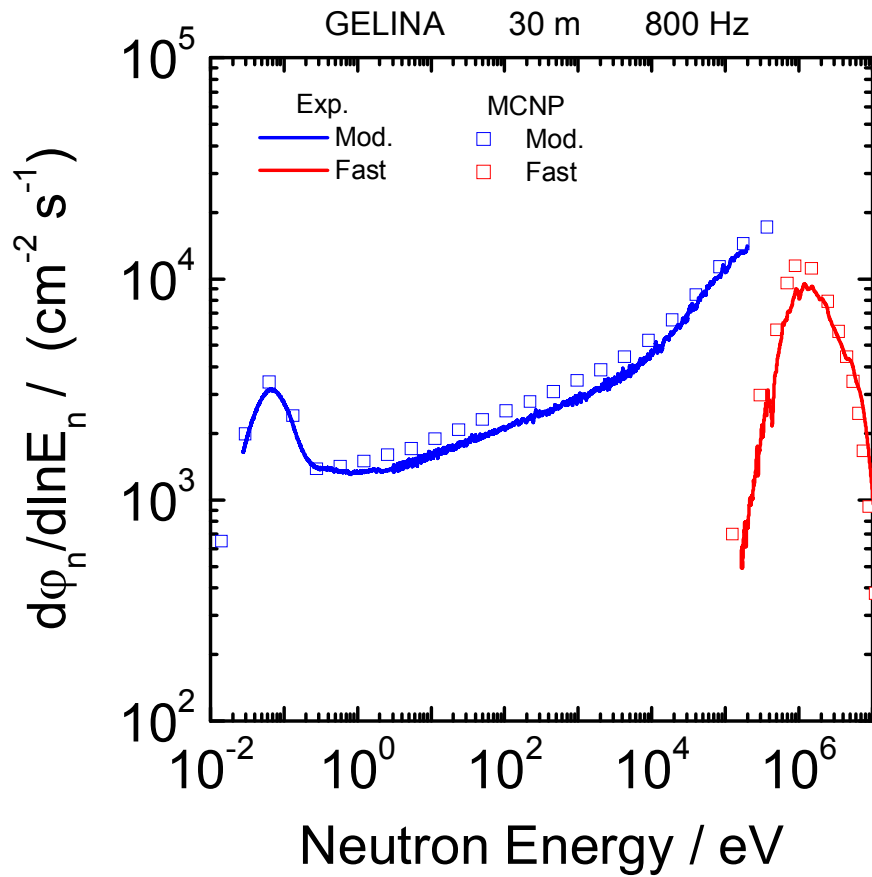
Flaska et al., NIM A 531, 394 (2004)

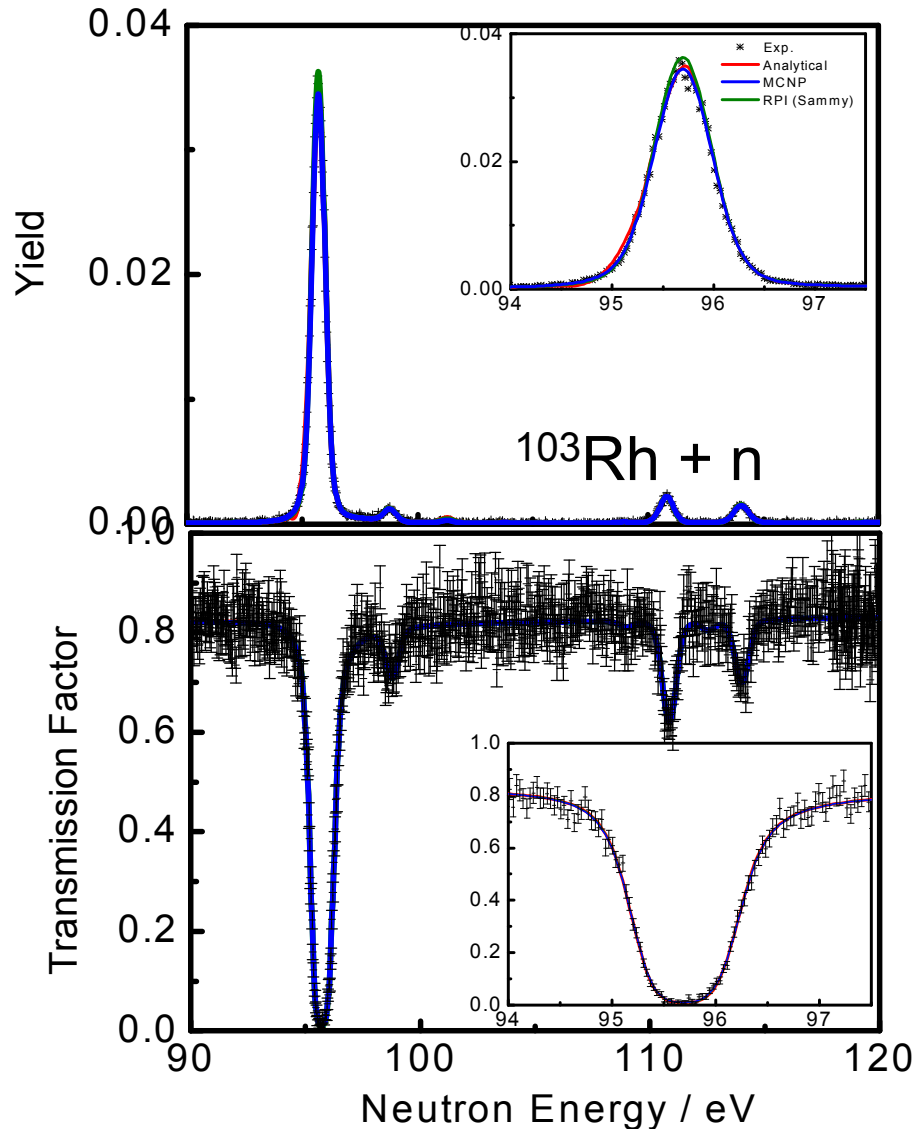
Flaska et al., NIM A 555, 329 (2005)

Neutron Flux



Resolution Function





Resolution Function
(RF)

$E_0 = 95.7 \text{ eV}$
 Γ_n / meV

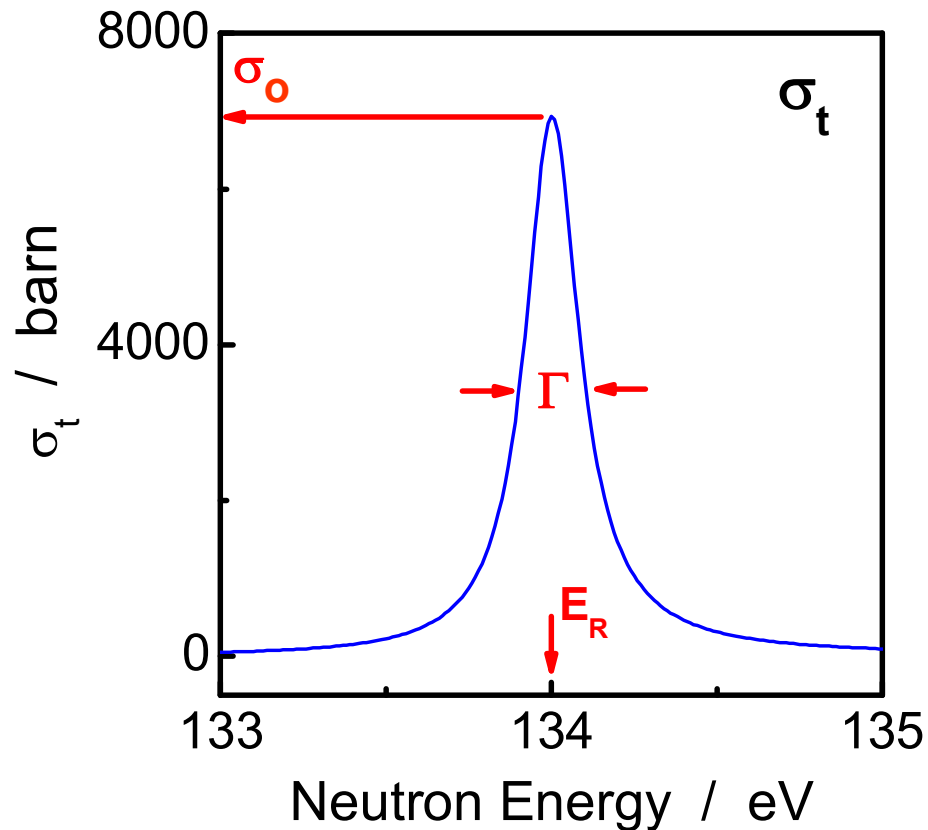
Capture

Analytical (REFIT)	2.47 (5)
MCNP	2.42 (5)
RPI (fit MCNP RF)	2.50 (5)

Transmission

Analytical (REFIT)	2.40 (5)
MCNP	2.42 (5)
RPI (fit MCNP RF)	2.40 (5)

- ⇒ RF has an impact on RP deduced from RSA
- ⇒ Transmission less sensitive to RF than capture



A cross section as a function of E_n shows a resonant structure, which can be described by a Breit-Wigner shape :

$$\sigma_t \sim \frac{1}{(E_n - E_R)^2 + (\Gamma/2)^2}$$

with

Γ natural line width (FWHM)

E_R resonance energy

- **Total Cross Section** (n,tot)

$$\sigma_t(E_n) = g_J \frac{\pi}{k_n^2} \frac{\Gamma_n \Gamma}{(E_n - E_R)^2 + (\Gamma/2)^2} \quad g_J = \frac{2J+1}{2(2I+1)}; k_n = \text{wavenumber}$$

- **Elastic Cross Section** (n,n)

$$\sigma_n(E_n) = \sigma_t(E_n) \frac{\Gamma_n}{\Gamma}$$

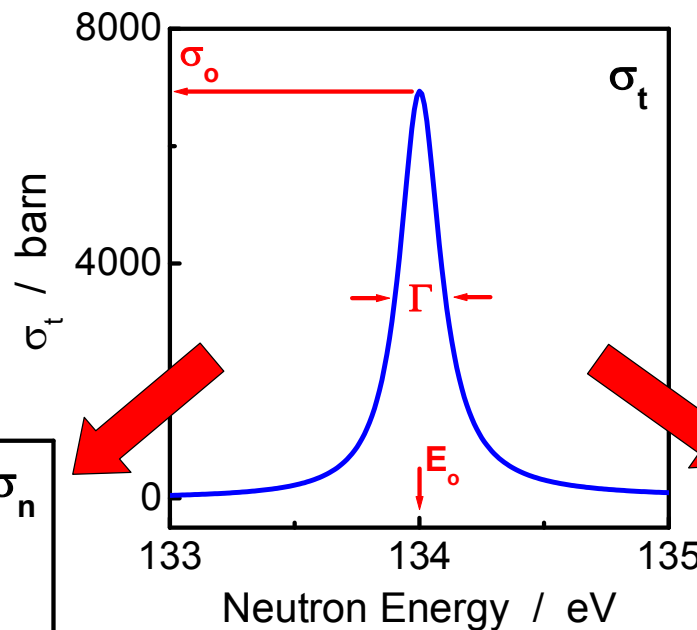
- **Capture Cross Section** (n, γ)

$$\sigma_\gamma(E_n) = \sigma_t(E_n) \frac{\Gamma_\gamma}{\Gamma}$$

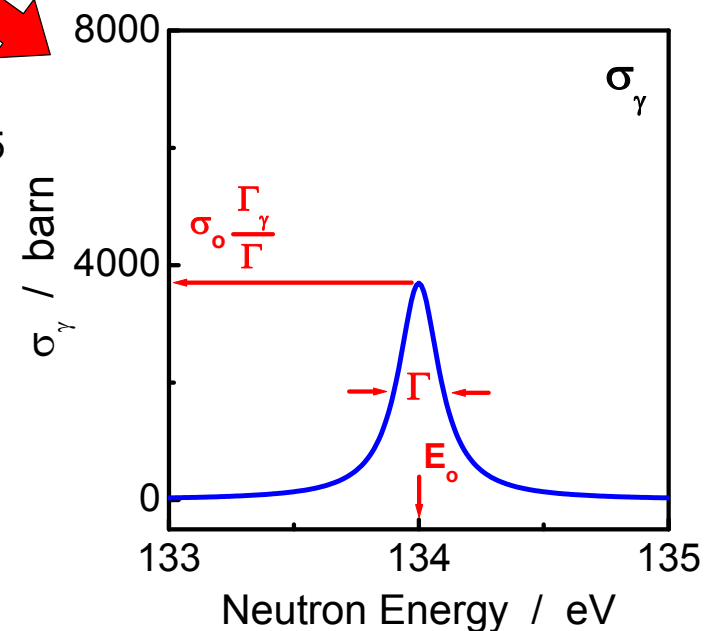
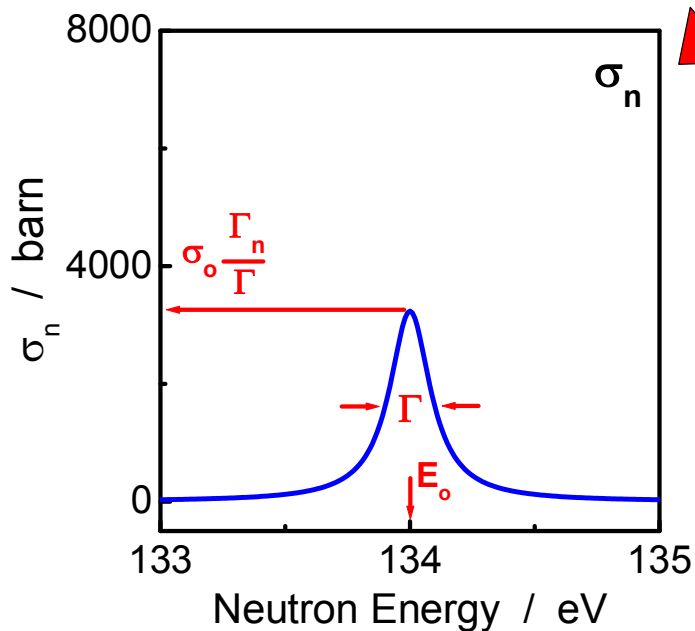
e.g. ^{109}Ag s-wave at $E_0 = 134$ eV

$(E_R, \Gamma_n, \Gamma_\gamma, J^\pi, l)$

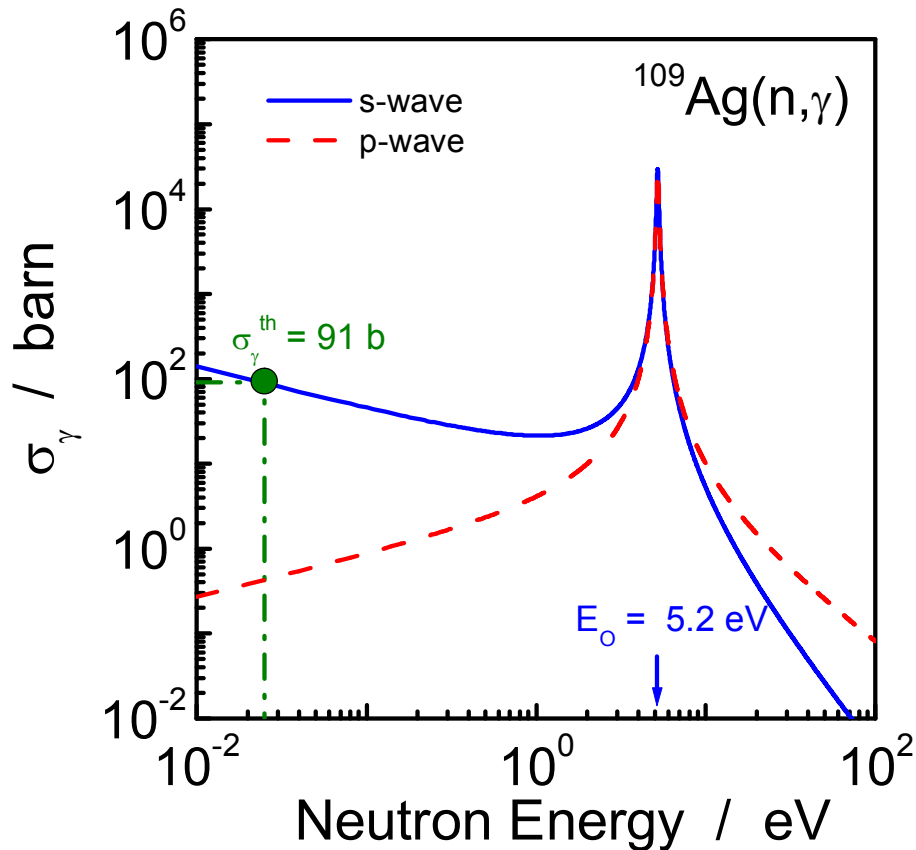
$$\sigma_0 = \frac{4\pi}{k_R^2} g_J \frac{\Gamma_n}{\Gamma}$$



$E_R = 134$ eV
 $\Gamma_n = 0.093$ eV
 $\Gamma_\gamma = 0.106$ eV
 $J^\pi = 1^-$
 $l = 0$
 $g_J = 3/4$



$$\sigma_{\gamma}(E_n) = g_J \frac{\pi}{k_n^2} \frac{\Gamma_n \Gamma_{\gamma}}{(E_n - E_R)^2 + (\Gamma/2)^2}$$



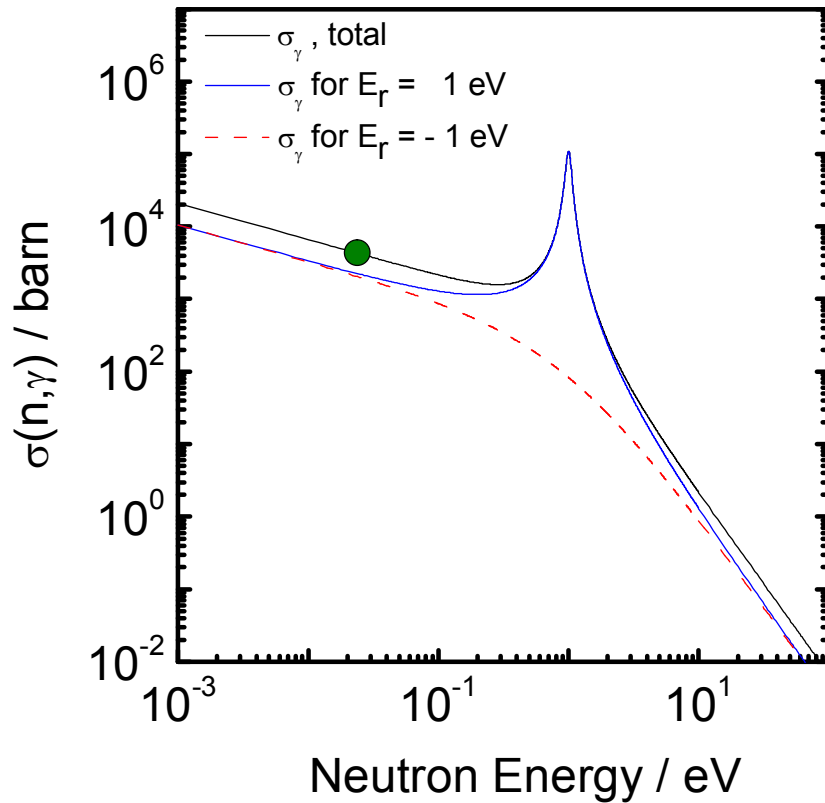
Γ_n depends on E_n : due to the centrifugal-barrier penetrability, which depends on the angular momentum of the incoming neutron l and E_n

(penetration probability depends on l)

- **s-wave** $\Gamma_n(E_n) = \Gamma_n^0 \sqrt{\frac{E_n}{1\text{eV}}}$
- **p-wave** $\Gamma_n(E_n) = \Gamma_n^1 \sqrt{\frac{E_n}{1\text{eV}}} \frac{k_n^2 a^2}{1 + k_n^2 a^2}$

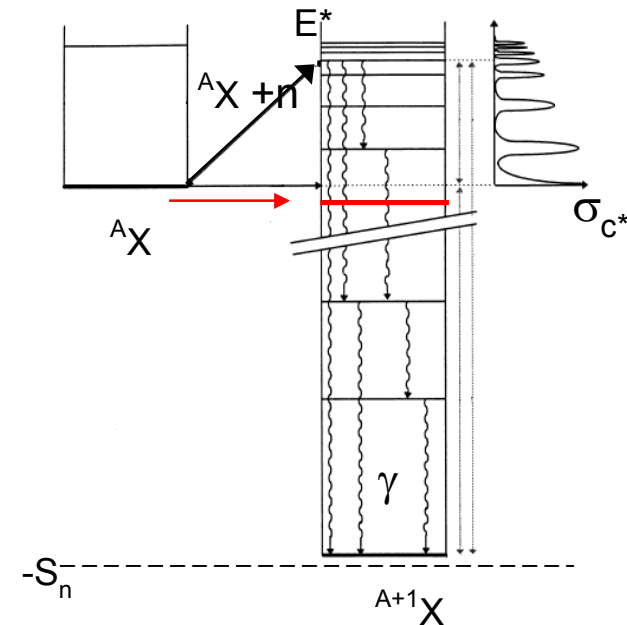
\Rightarrow **cross section at 0.025 eV (Thermal)**

$$\sigma_r^{\text{th}} \approx 4.099 \cdot 10^6 \left(\frac{A+1}{A} \right)^2 \sum_{j=1}^N \frac{g_J \Gamma_{nj}^0 \Gamma_{rj}}{E_{Rj}^2}$$



$$\bullet \sigma^{th} > \sigma_{E>0}^{th} \cong 4.099 \cdot 10^6 \left(\frac{m_A + 1}{m_A} \right)^2 \sum_{j=1}^N \frac{g_j \Gamma_{nj}^0 \Gamma_{\gamma j}}{E_{Rj}^2}$$

• Additional contribution from bound states (negative resonances)



- (n,γ)

$$\sigma_{\gamma}(E_n) = g_J \frac{\pi}{k_n^2} \frac{\Gamma_n \Gamma_{\gamma}}{(E_n - E_R)^2 + (\Gamma/2)^2}$$

$$\Gamma = \Gamma_n + \Gamma_{\gamma}$$

$$g_J = \frac{2J+1}{2(2I+1)}$$

$$R = (0.80 + 1.23A^{1/3}) \text{ fm}$$

- (n,n)

$$\sigma_n(E_n) = g_J \frac{\pi}{k_n^2} \frac{\Gamma_n \Gamma_n}{(E_n - E_R)^2 + (\Gamma/2)^2} + g_J \frac{4\pi}{k_n} \frac{\Gamma_n (E - E_R) R}{(E_n - E_R)^2 + (\Gamma/2)^2} + g_J 4\pi R^2$$

- Total

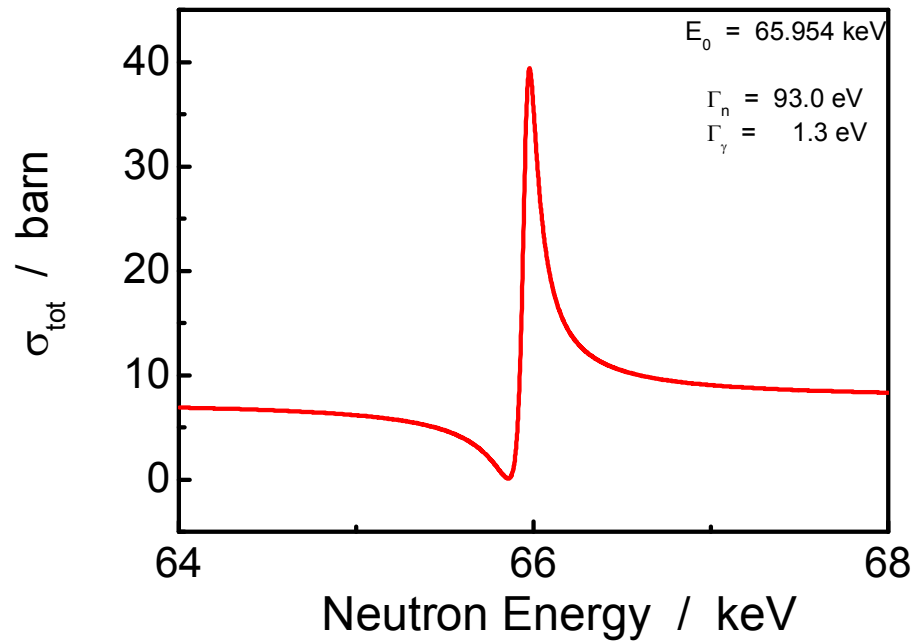
$$\sigma_{\text{tot}}(E_n) = g_J \frac{\pi}{k_n^2} \frac{\Gamma_n \Gamma}{(E_n - E_R)^2 + (\Gamma/2)^2} + g_J \frac{4\pi}{k_n} \frac{\Gamma_n (E - E_0) R}{(E_n - E_R)^2 + (\Gamma/2)^2} + g_J 4\pi R^2$$

Interference term

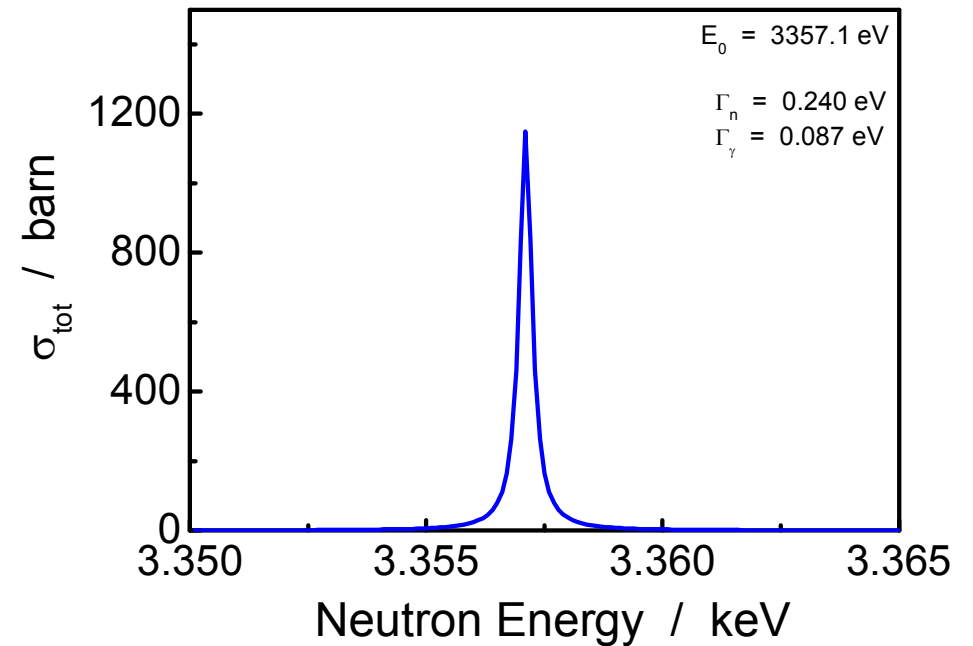
s-wave <----> p-wave ($l > 0$)

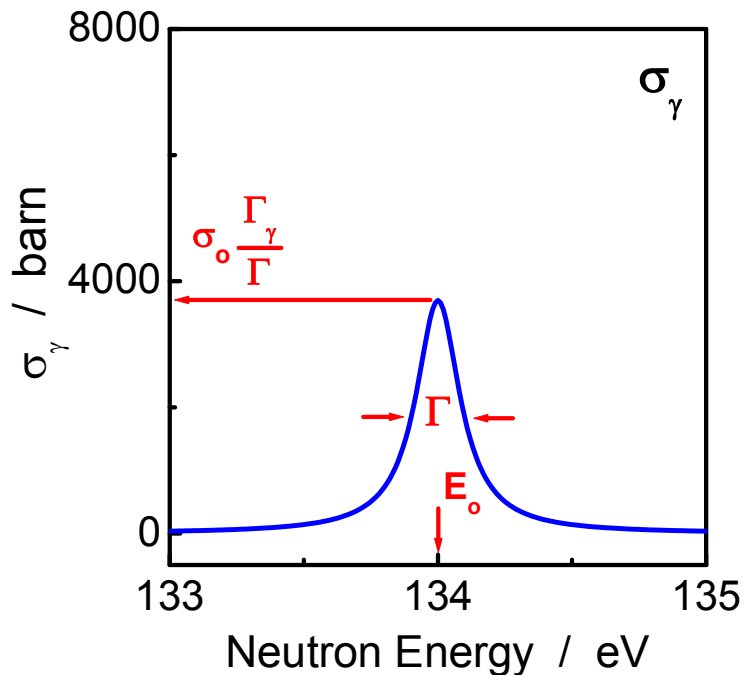
$$\sigma_{\text{tot}}(E_n) = g_J \frac{\pi}{k_n^2} \frac{\Gamma_n \Gamma}{(E_n - E_R)^2 + (\Gamma/2)^2} + g_J \frac{4\pi}{k_n} \frac{\Gamma_n (E - E_0) R}{(E_n - E_R)^2 + (\Gamma/2)^2} - g_J 4\pi a^2$$

s-wave



p-wave





- In **ideal conditions**, without any instrumental resolution broadening, we can determine $\sigma_0 \Gamma_\gamma / \Gamma$ and Γ from one experiment.

$$\Rightarrow \Gamma_\gamma \text{ and } \Gamma_n$$

$$\sigma_0 = f(g_J, \Gamma_n, \Gamma)$$

However,

- Due to instrumental limitations it is mostly impossible to determine Γ and $\sigma_0 \Gamma_\gamma / \Gamma$
- **Due to the instrumental limitations** the effective experimental observable is mostly:
the area of a resonance

- Transmission**

$$A_{t,thin} \propto ng_J \Gamma_n$$

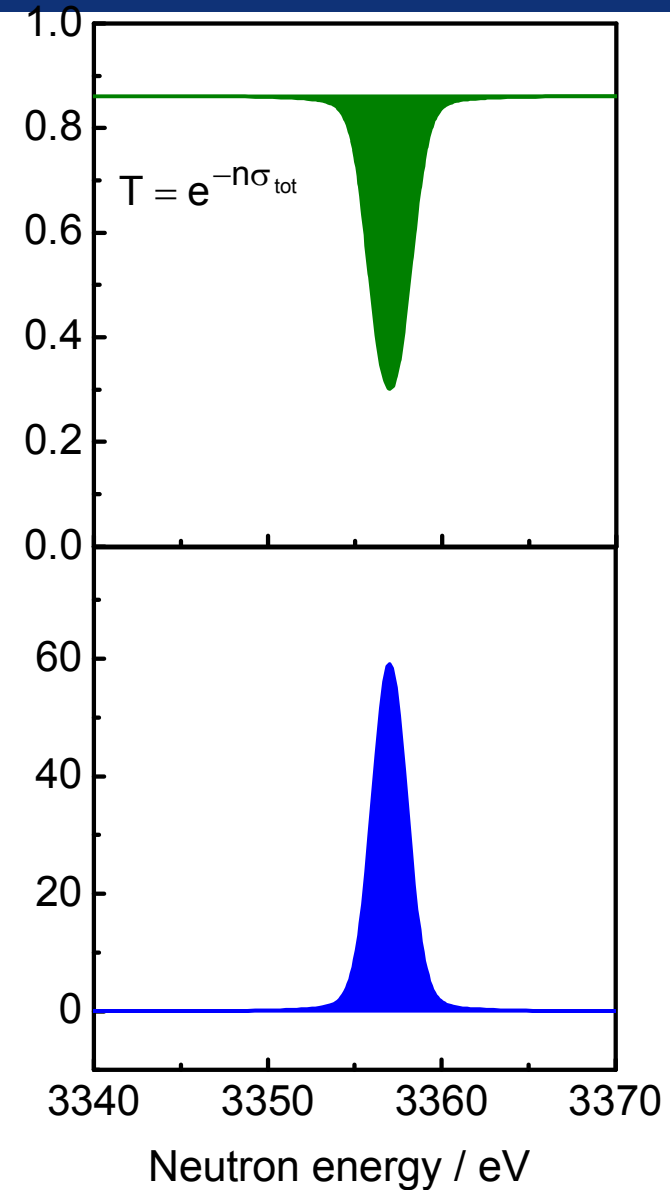
$$A_{t,thick} \propto \sqrt{ng_J \Gamma_n \Gamma}$$

- Capture**

$$A_{\gamma,thin} \propto ng_J \frac{\Gamma_n \Gamma_\gamma}{\Gamma}$$

n is the target thickness in atoms/barn

Rae et al., Nucl. Phys. 5 (1958) 89
 F. Fröhner et al., ND1966, p. 55



Rae et al., Nucl. Phys. 5 (1958) 89
F. Fröhner et al., ND1966, p. 55

$$\Gamma_\gamma \ll \Gamma_n$$

$$\Gamma_\gamma \gg \Gamma_n$$

• Capture	$A_\gamma \propto n g_J \frac{\Gamma_n \Gamma_\gamma}{\Gamma}$	$n g_J \underline{\Gamma_\gamma}$	$n g_J \underline{\Gamma_n}$
• Transmission	$A_{t, \text{thin}}$	$n g_J \underline{\Gamma_n}$	$n g_J \underline{\Gamma_n}$
	$A_{t, \text{thick}}$	$\sqrt{n g_J \Gamma_n}$	$\sqrt{n g_J \Gamma_n \Gamma_\gamma}$

⇒ Ideally : **combine** thin capture measurements with transmission measurements on samples with different thicknesses

- **SLBW and Resonance Area Analysis**

- Instructive

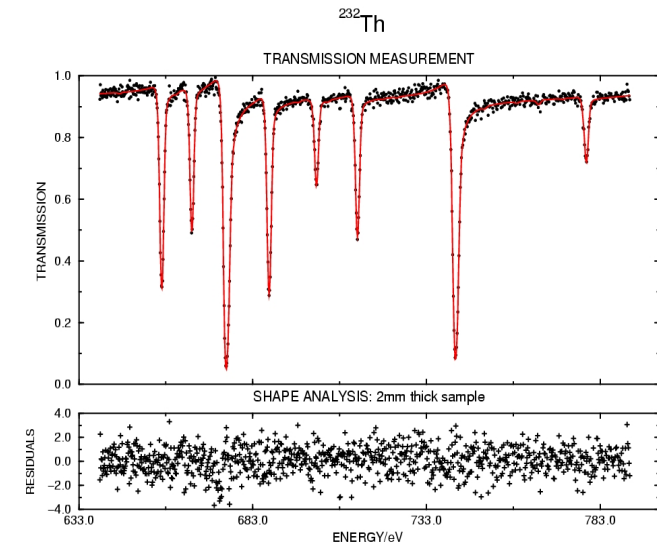
- **Resonance Shape Analysis (RSA)**

(Reich-Moore (or better) Recommended)

- Extension of RRR
- Direct treatment of broadening effects (Resolution & Doppler)
- Direct treatment of multiple scattering effects (capture data)
- Reduction of uncertainties

» **SAMMY (N. Larson)**

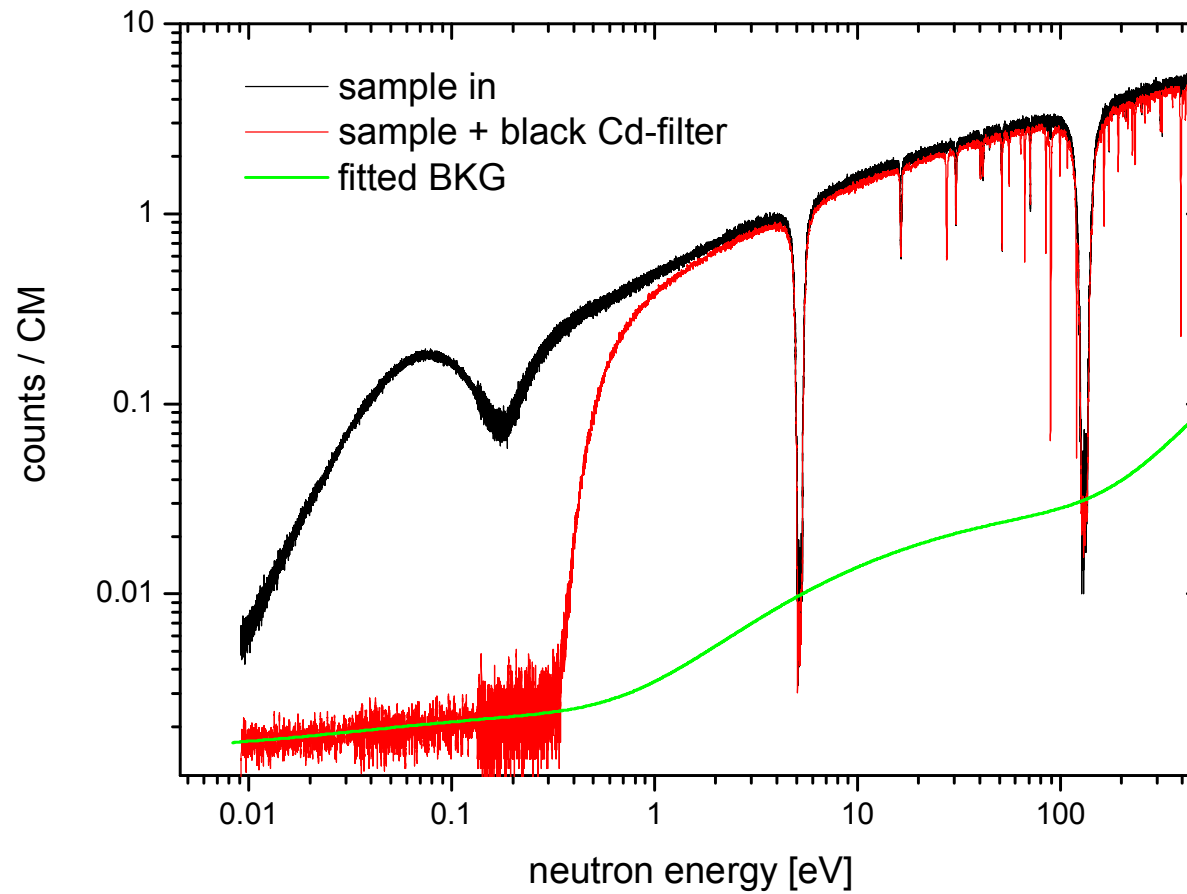
» **REFIT (M. Moxon)**

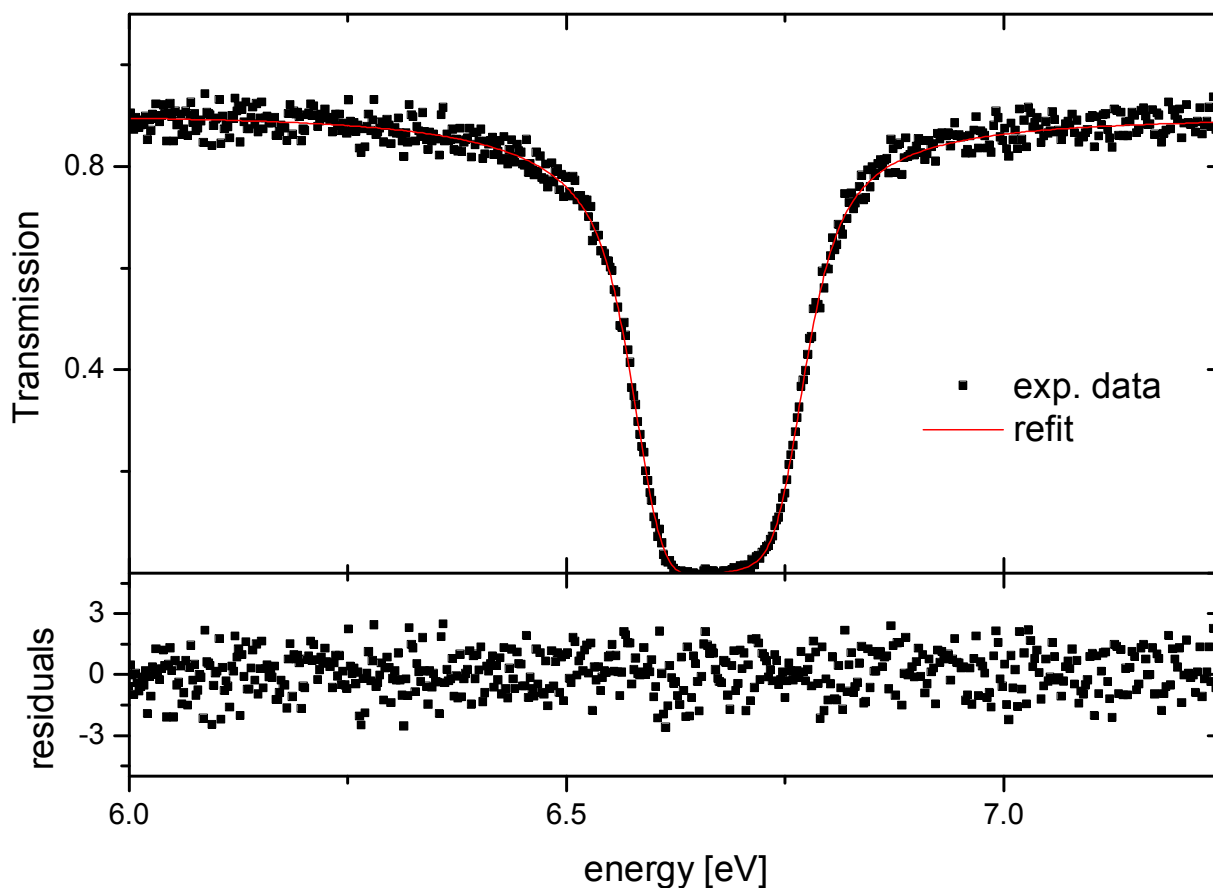


- **Cadmium**

	sample	energy	Γ_{γ} [meV]	Γ_n [meV]
Rainwater	metal	0.176+/-0.002	115 +/- 2	0.620 +/- 0.020
Brockhouse	metal	0.180 +/- 0.003	113 +/- 2	0.680 +/- 0.020
Akyuz	metal	0.181 +/- 0.003	109 +/- 3	0.645 +/- 0.025
Widder	powder	0.1776 +/- 0.0006	114.3 +/- 0.6	0.618 +/- 0.003
Harz	various	0.1783 +/- 0.0002	113.5 +/- 0.5	0.650 +/- 0.005

thickness (atoms/barn)	type	flight path length
1.38 10 ⁻⁴	capture	10 m
3.40 10 ⁻⁴	capture	10 m
5.40 10 ⁻⁴	transmission	50 m
1.10 10 ⁻³	capture	30 m
2.36 10 ⁻³	capture	30 m
4.67 10 ⁻³	capture	30 m
9.34 10 ⁻³	capture	10 m
	transmission	25 m, 50 m
2.34 10 ⁻²	transmission	50 m
1.20 10 ⁻¹	transmission	25m, 50m



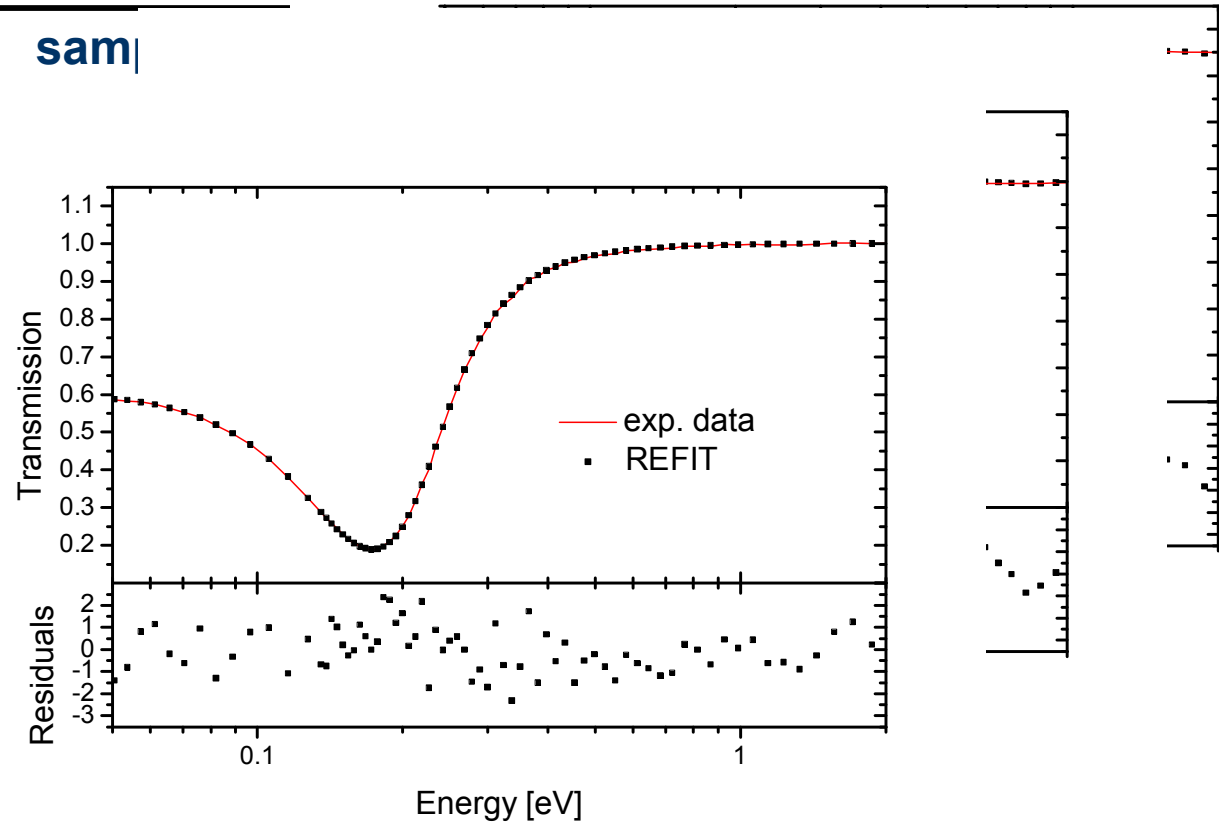


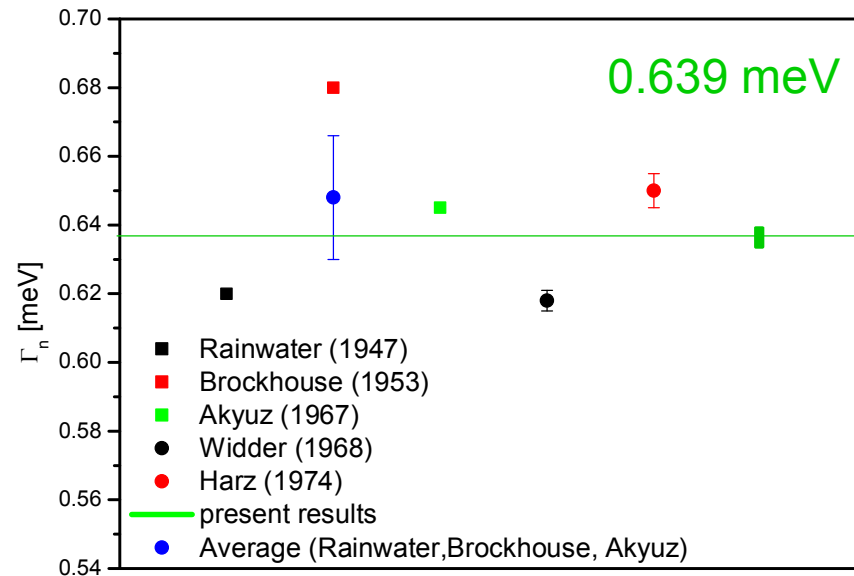
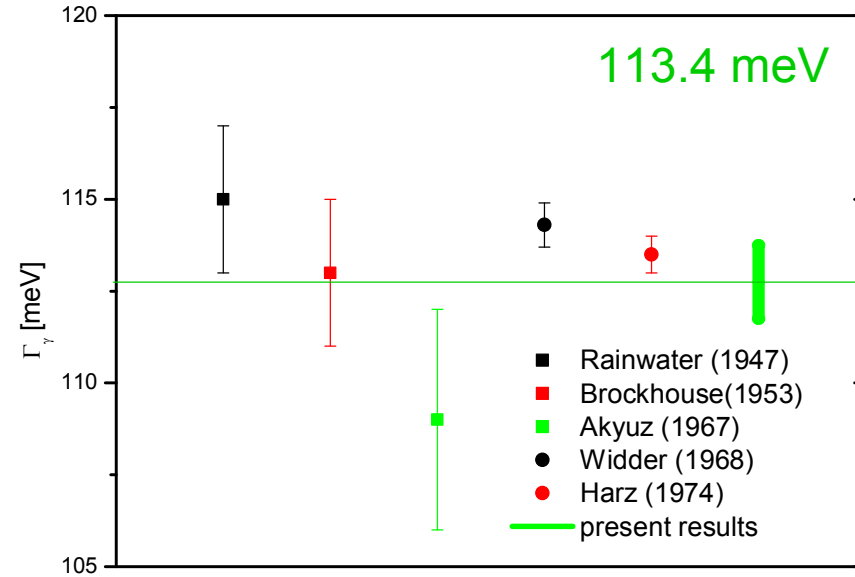
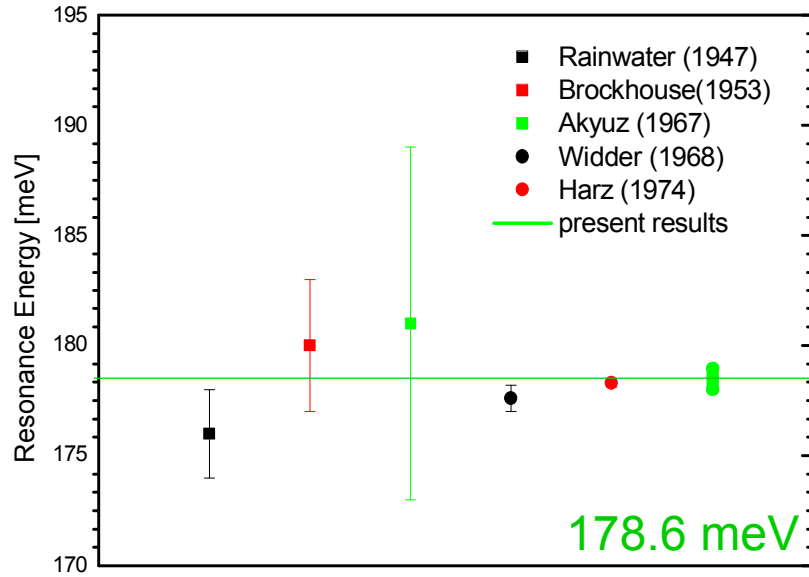
$$E_{\text{res}} = 6.6735 \pm 0.002$$

$$F_p = 26.4440 \pm 0.005$$

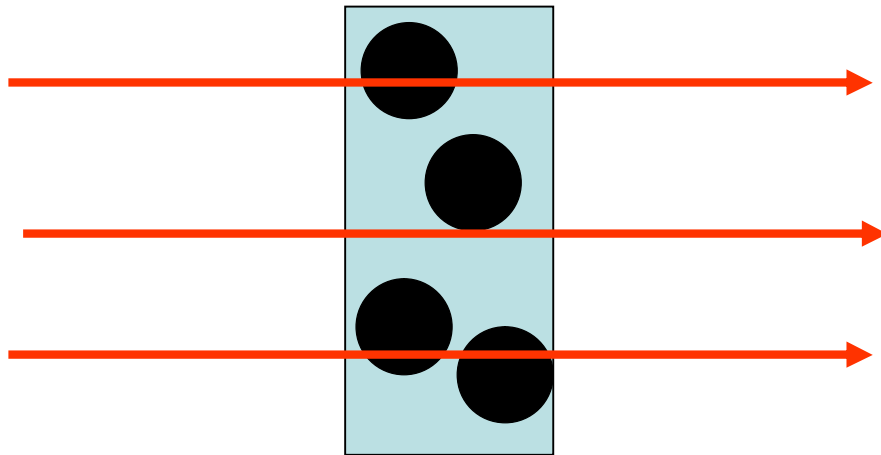
Checked consistency with
energies above 1 keV
(Oak Ridge, 200 m, data
by J. Harvey)

thickness (atoms/barn)
$1.40 \cdot 10^{-4}$
$1.36 \cdot 10^{-4}$
$2.24 \cdot 10^{-4}$
$2.80 \cdot 10^{-4}$

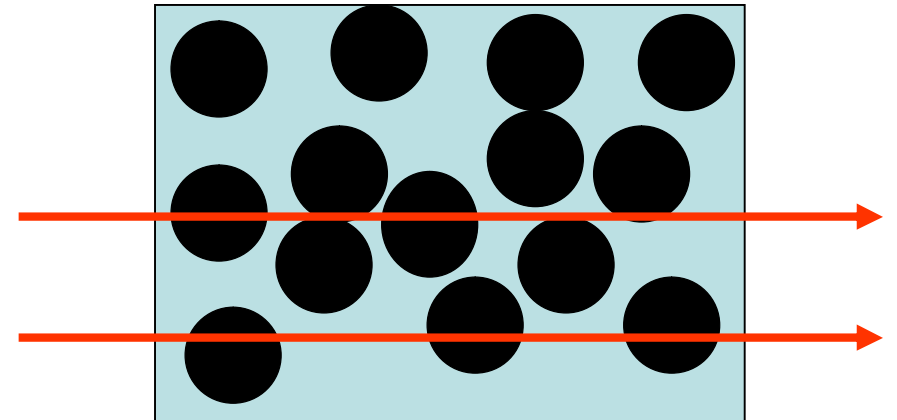




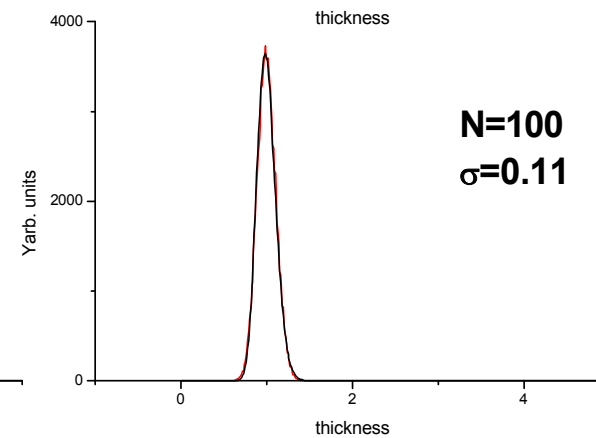
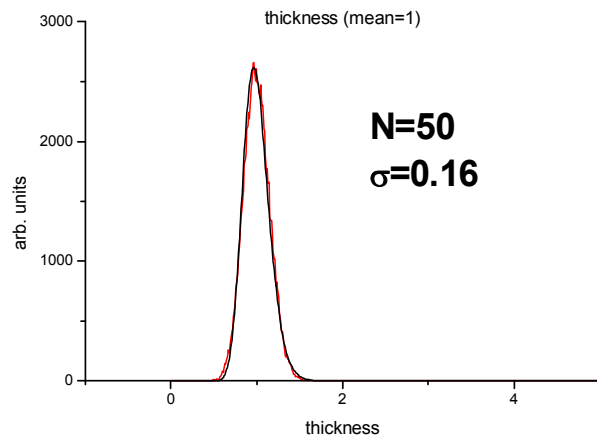
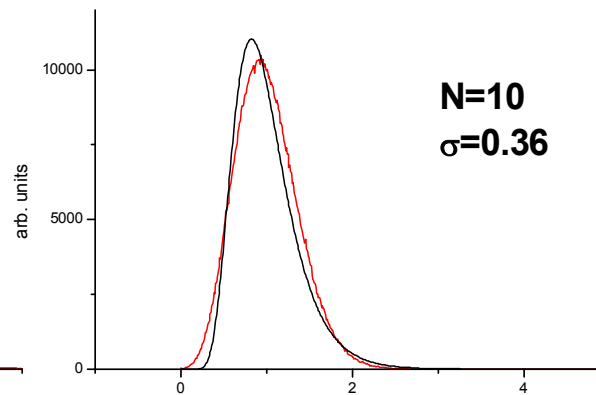
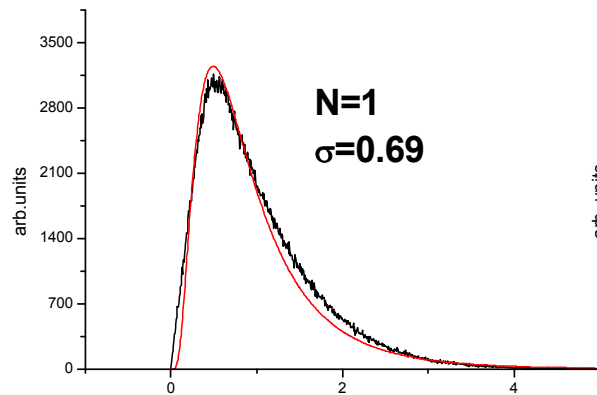
N=1



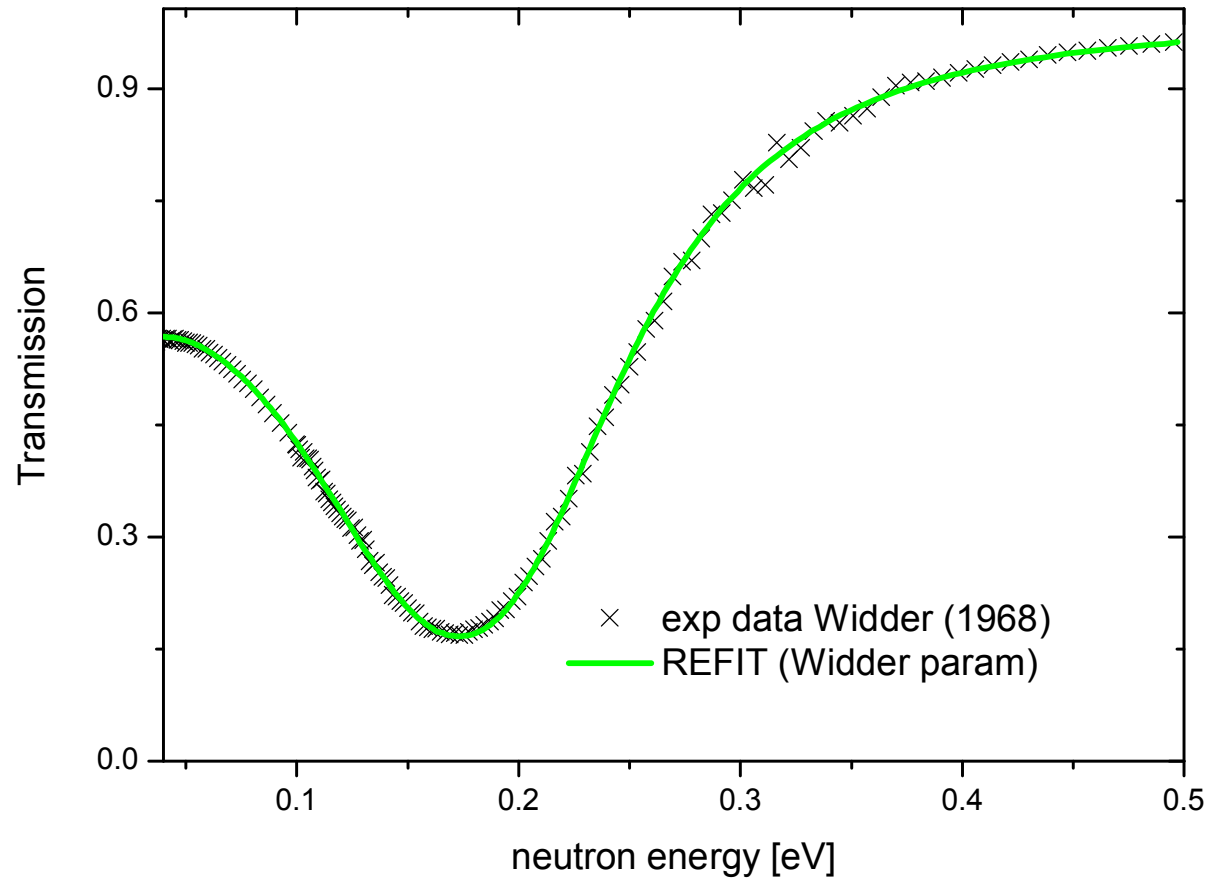
N=3.5



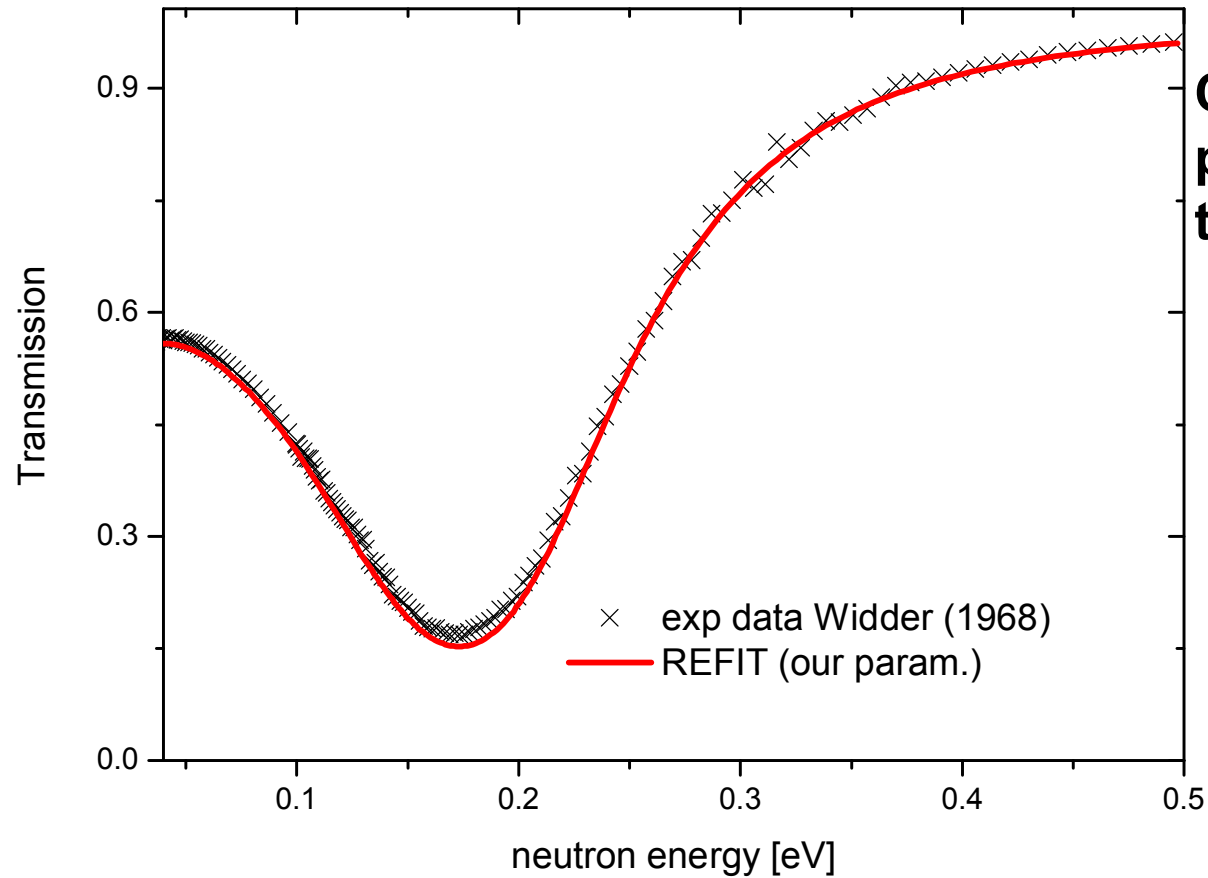
**Probability that a neutron “sees” n particle,
given by Poisson-statistics**



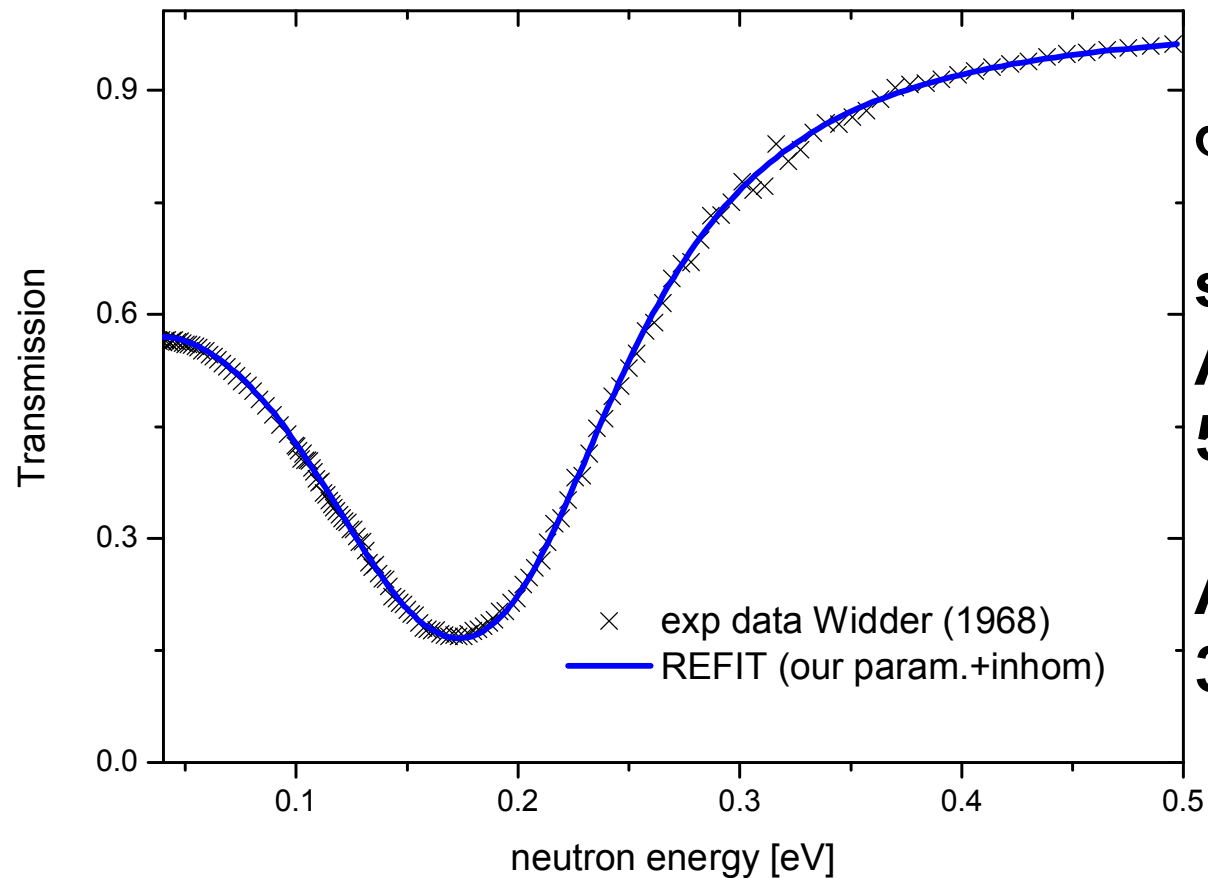
- Poisson distribution
- Spheres
- Radius lognormal
- Black line MC
- Red line lognormal fit



**Calculation using
parameters given by
Widder et al.**



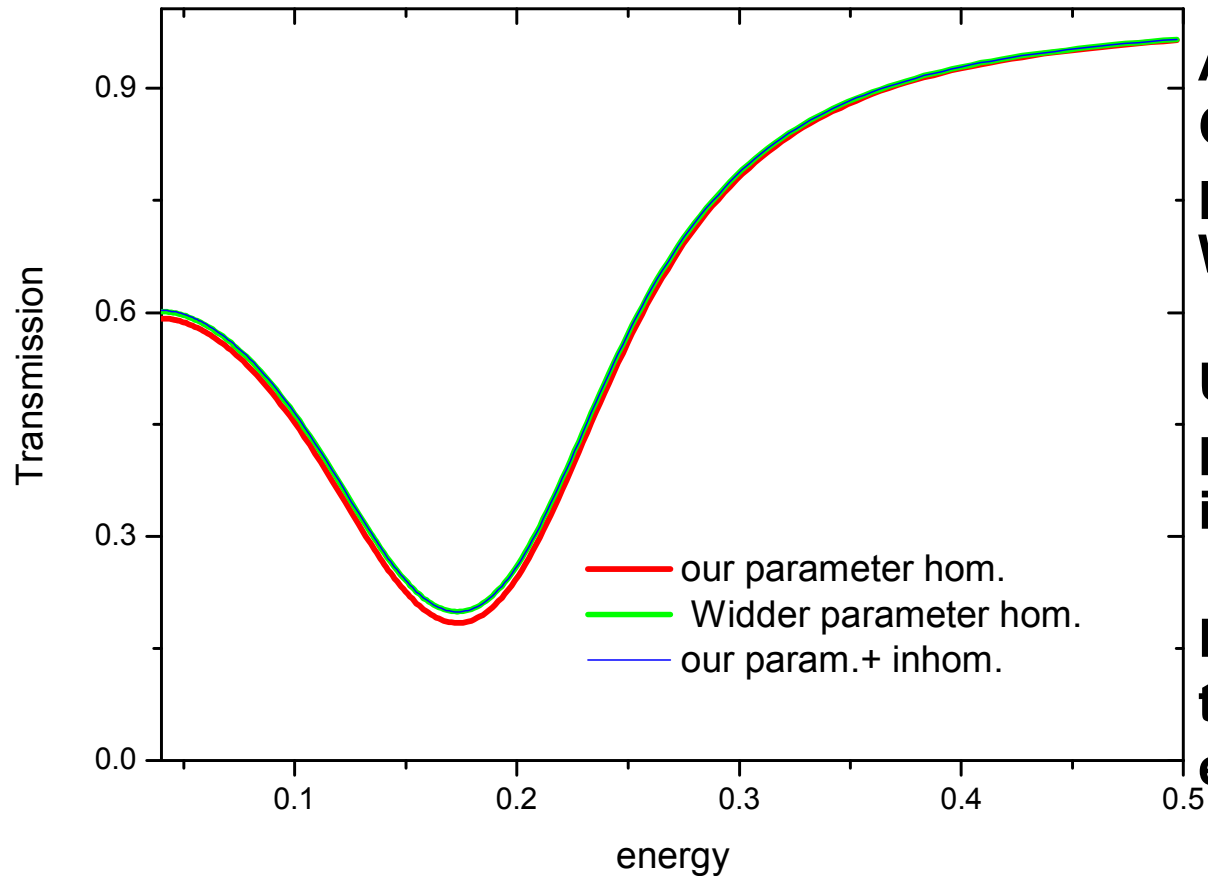
**Calculation using
parameters derived by fit
to “our data”**



$$\sigma_{\text{inhom}} = 0.096$$

sample thickness:
Approx. 50 μm
50-200 layers (estim)

Average particle size
300nm – 1 μm (estim)



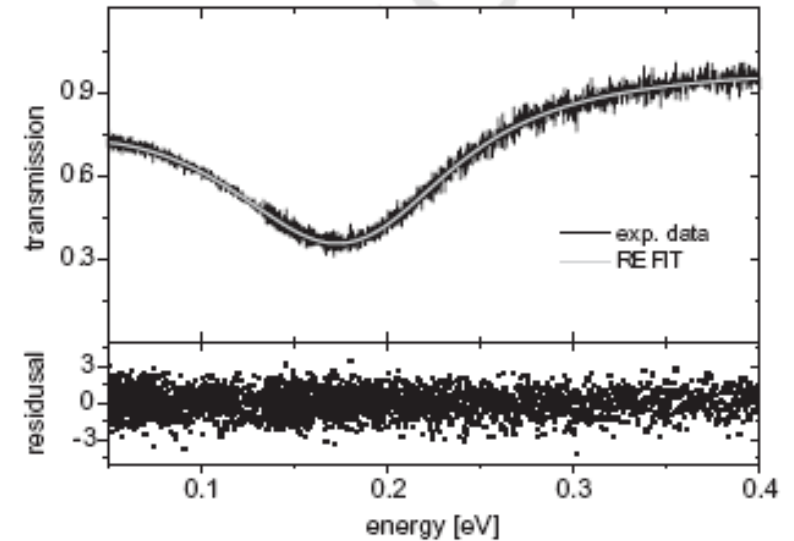
**Agreement between:
Calculation using
parameters given by
Widder et al.**

**Using *our* resonance
parameters+ sample
inhogeneity**

**Relative deviation less
than 0.001 over whole
energy range**

Conclusion: Widder et al. did a good experiment but their sample was not homogenous.

thickness (at/b)	sample type	flight path length
$1.40 \cdot 10^{-4}$	solution	25 m
$1.36 \cdot 10^{-4}$	foil	25m
$2.24 \cdot 10^{-4}$	foil	25 m

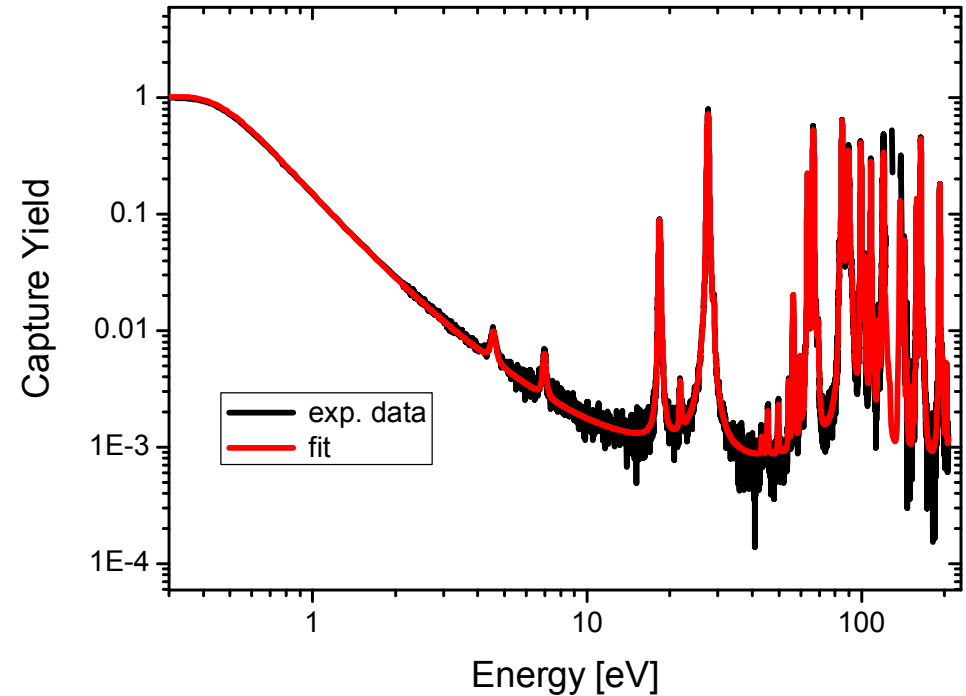
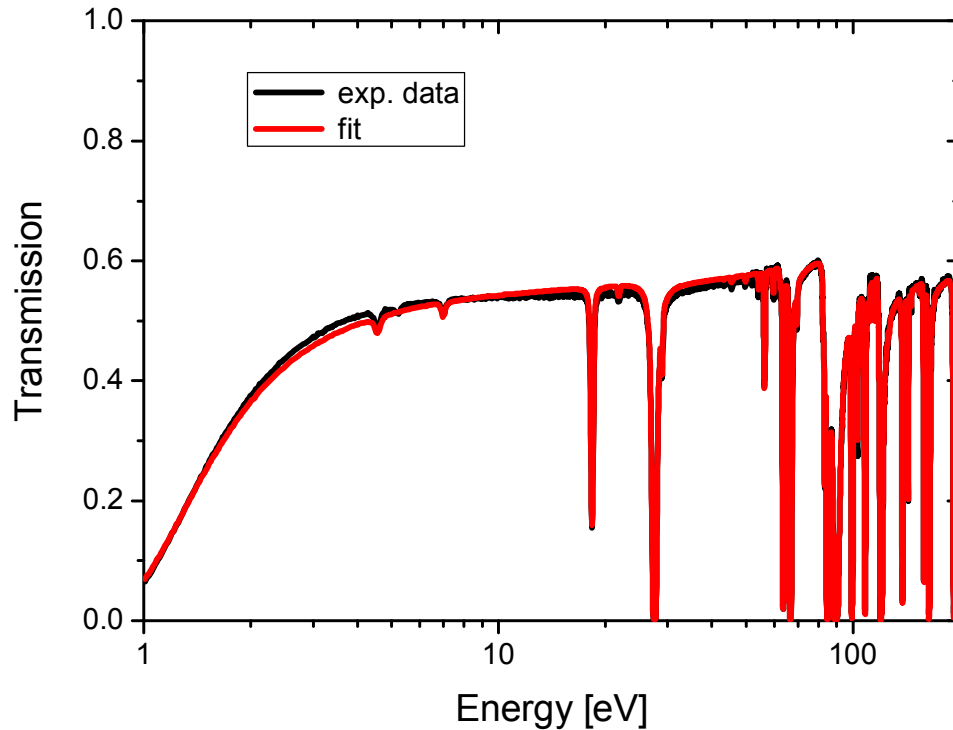


Only uncorrelated uncertainties

Parameter	p / meV	$\rho(p_i, p_j)$
E_R	178.7 ± 0.1	1.00 0.43 0.79
Γ_n	0.640 ± 0.001	1.00 0.43
Γ_γ	113.5 ± 0.2	1.00

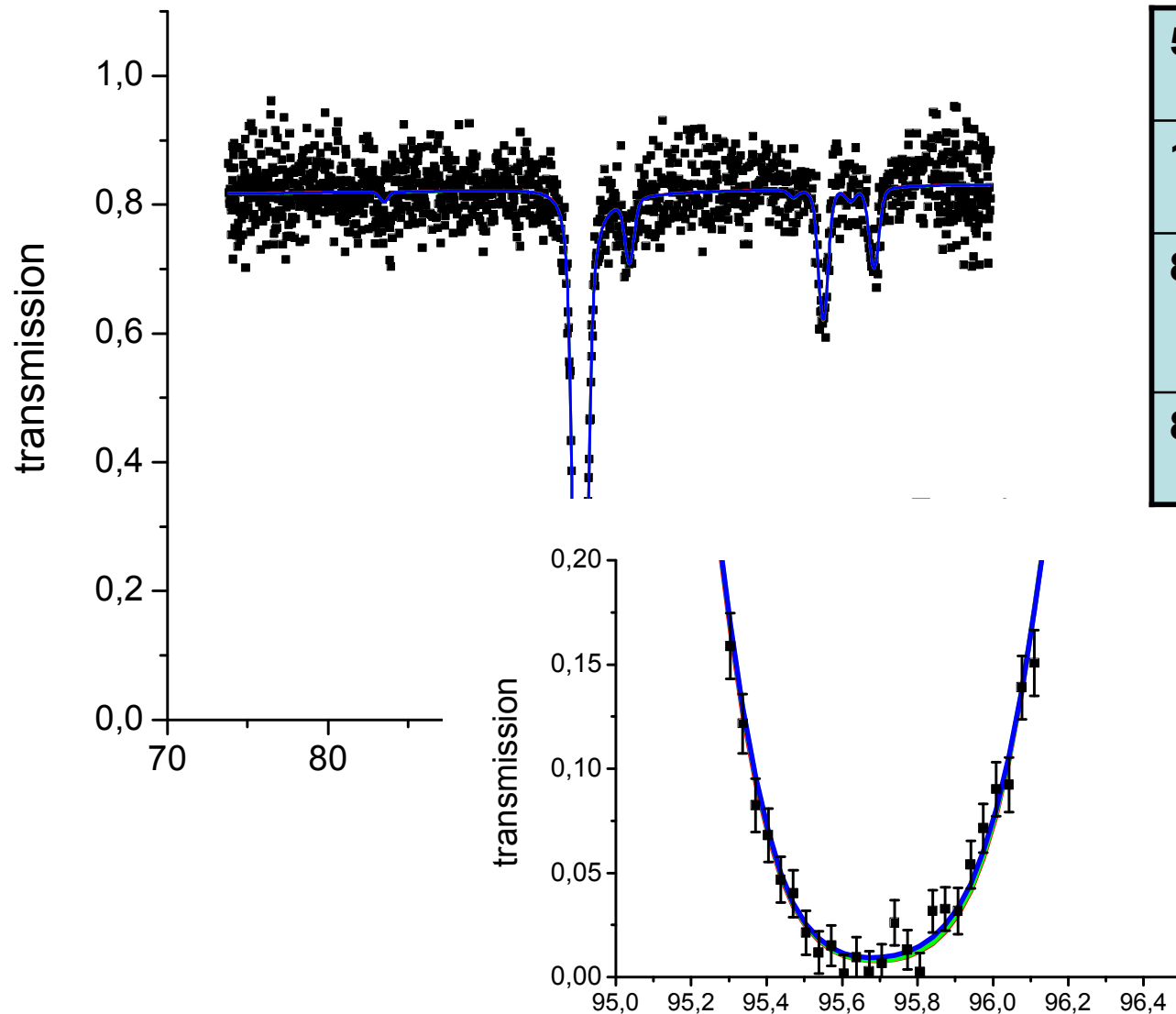
Full Covariance Matrix

Parameter	p / meV	$\rho(p_i, p_j)$
E_R	178.7 ± 0.1	1.00 0.53 0.28
Γ_n	0.640 ± 0.004	1.00 0.25
Γ_γ	113.5 ± 0.2	1.00

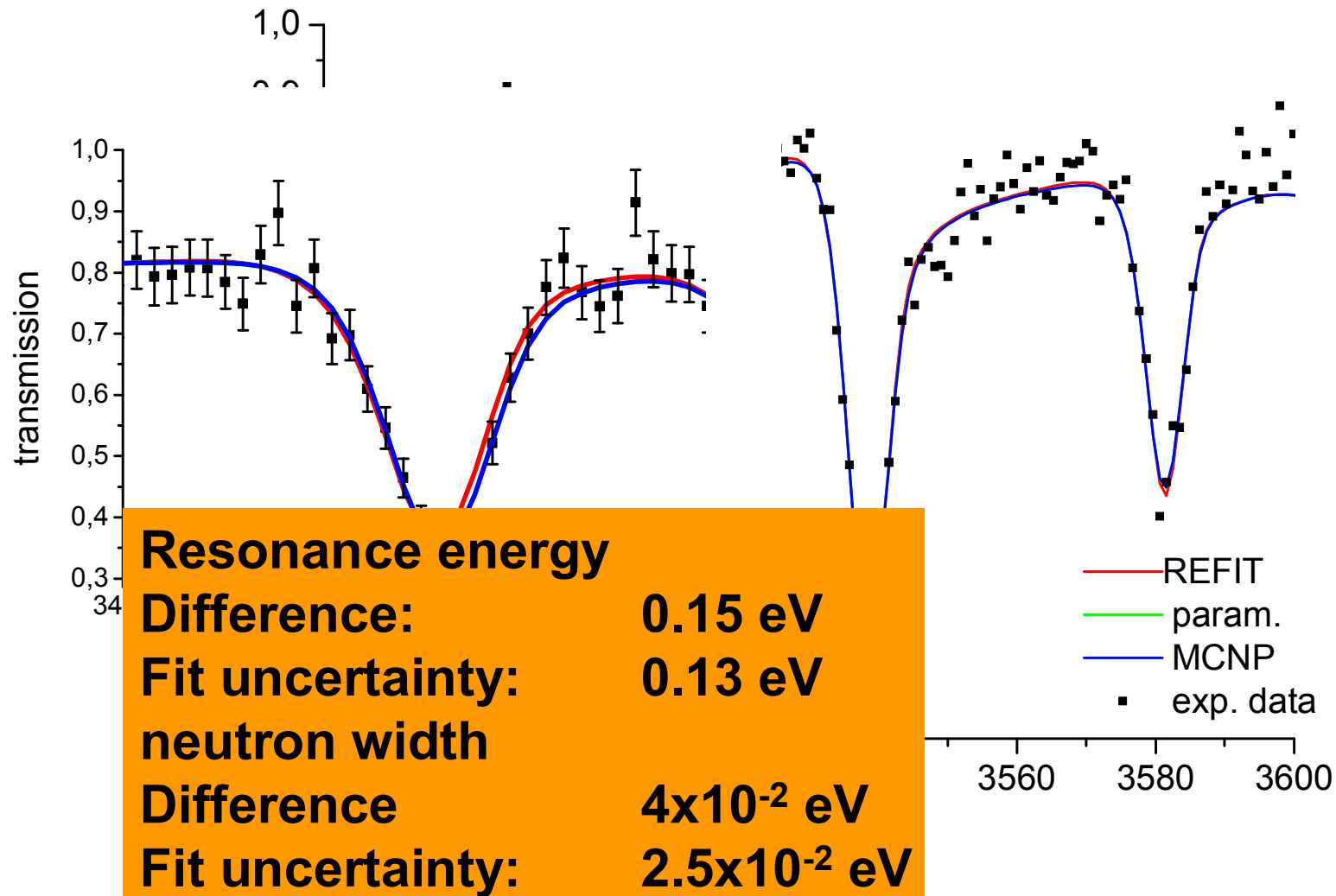


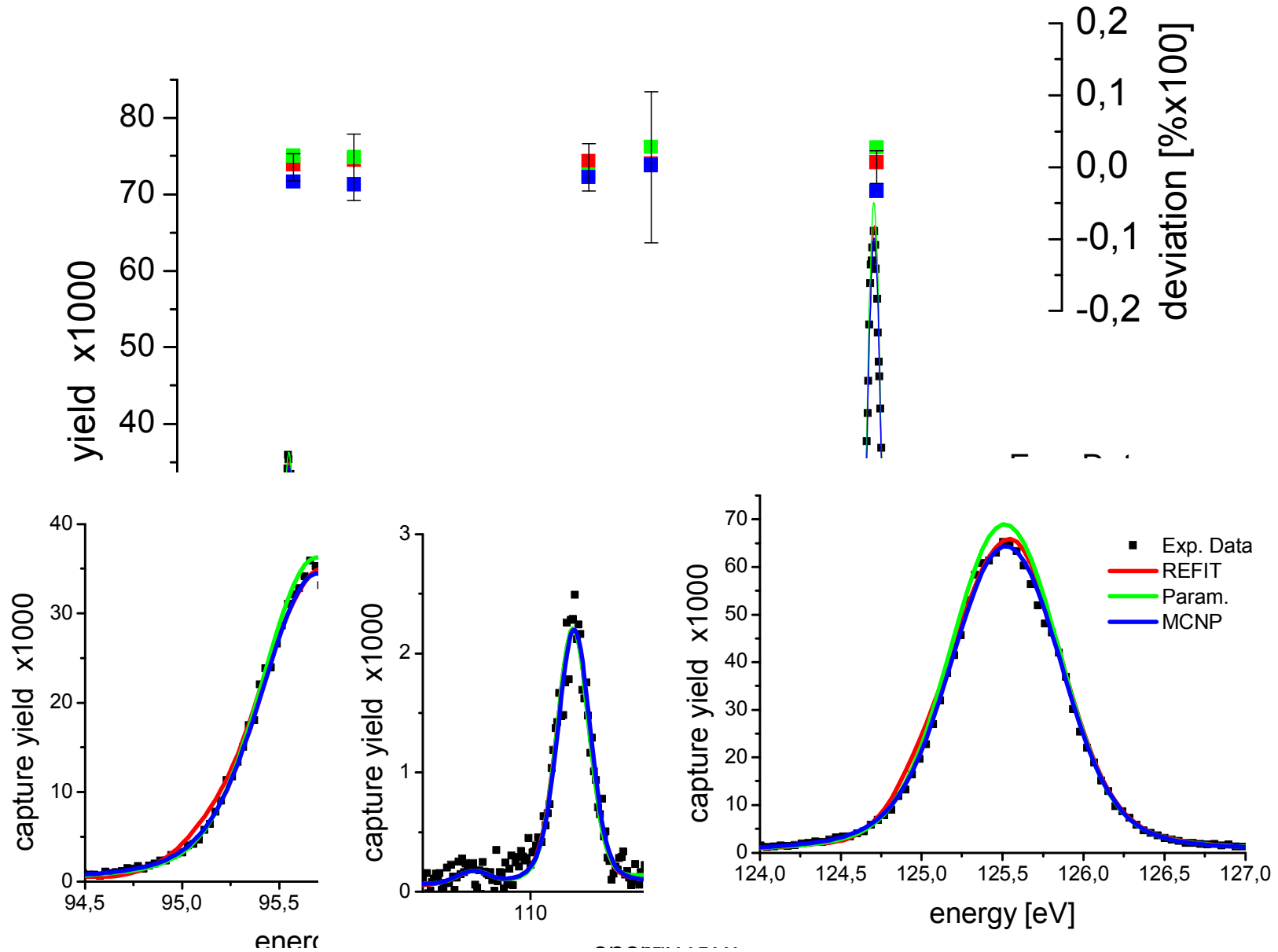
- capture/transmission RSA analysis finalized up to 1.5 keV
- input to the IRDF 2009 (IAEA) up to 1 keV
- consistency check including comparisons with results on enriched samples
- Thermal point difficult since Cd is not monoisotopic (but dominated by ^{113}Cd)
- comparisons with integral experiments (impact of first resonance parameters) (see NIMB 267)

• Rhodium



50 Hz	10 m	thick	capture
100 Hz	50 m	thin	transm.
800 Hz	30 m	thin/	capture
	60 m	thick	
800 Hz	50 m	very thick	transm.



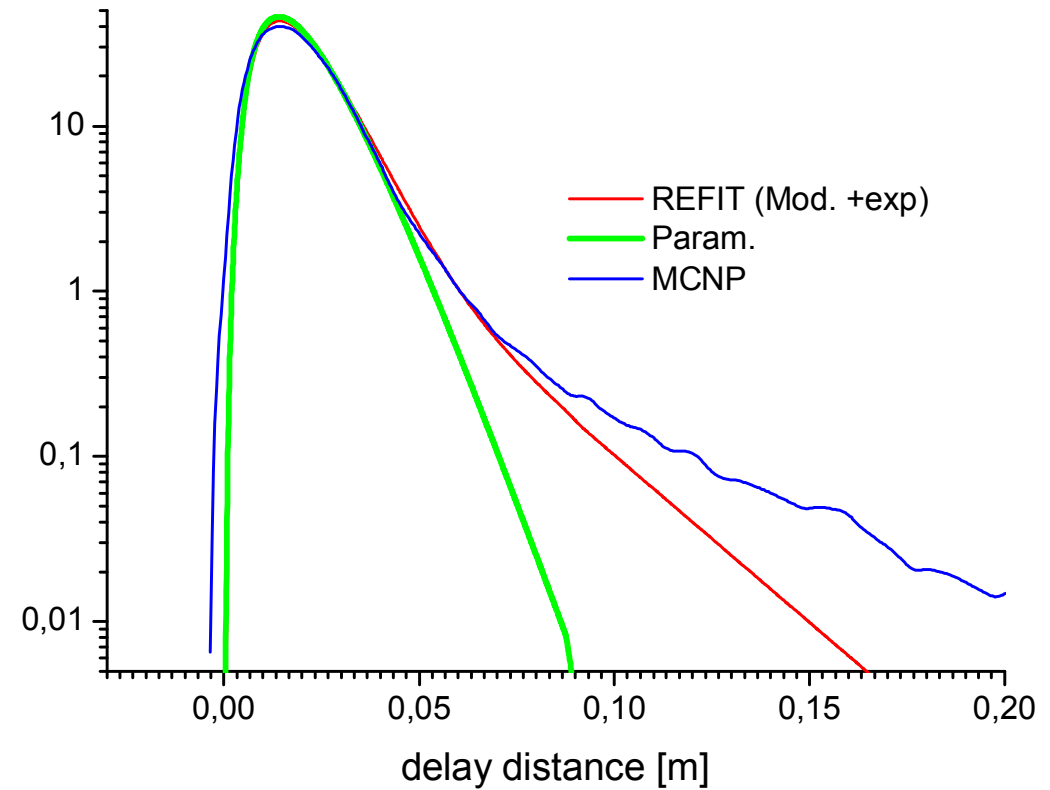


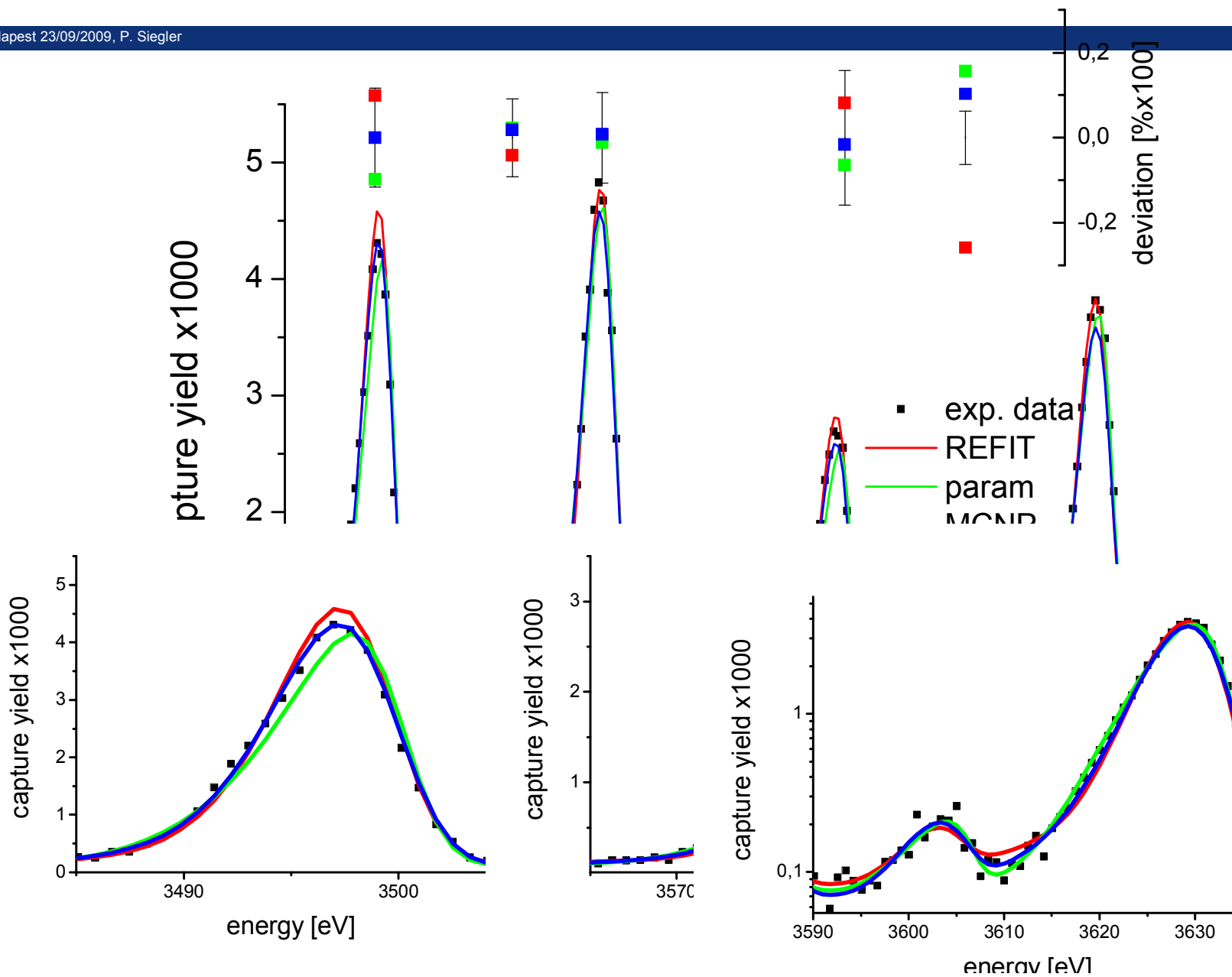
- REFIT

- Exponential Decay

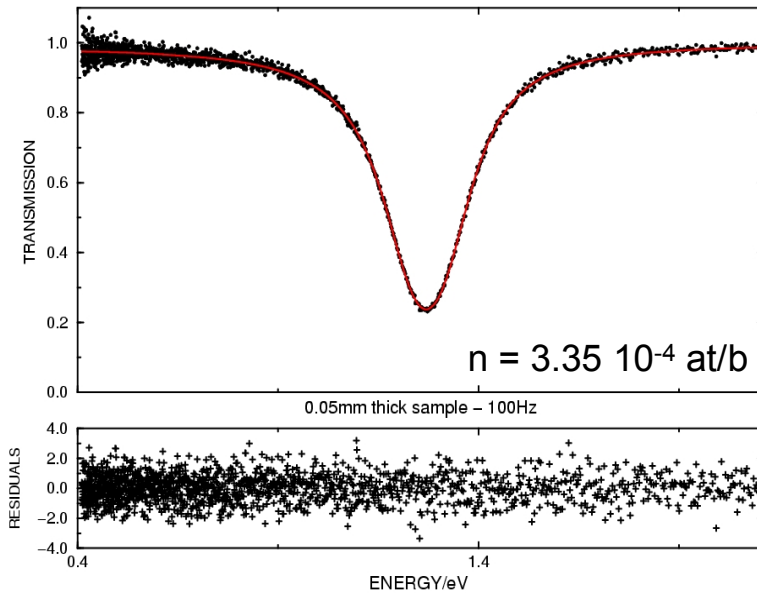
- Amplitude 8.5 %
- Decay constant 110 ns

Prob. density





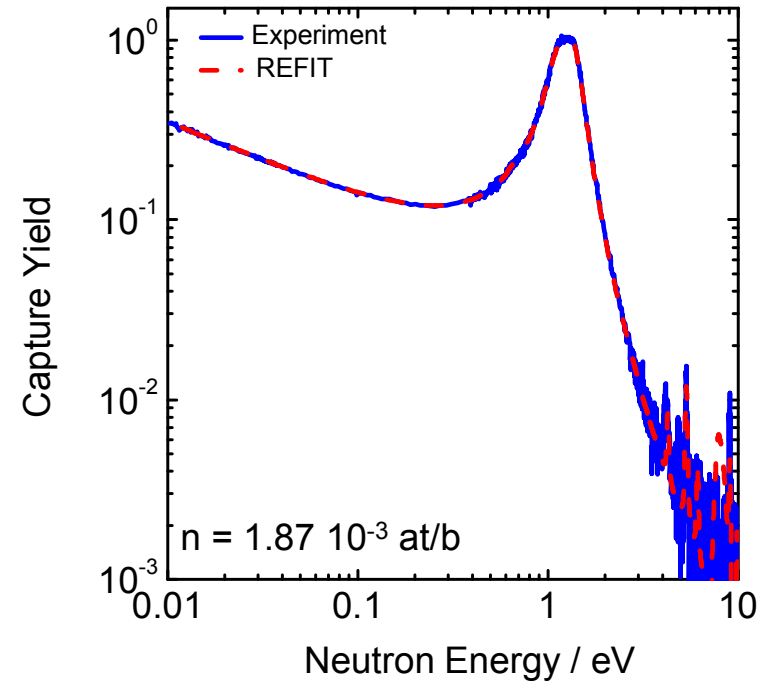
Transmission: 50 m 100 Hz



REFIT
RP + N



Capture : 15 m 40 Hz



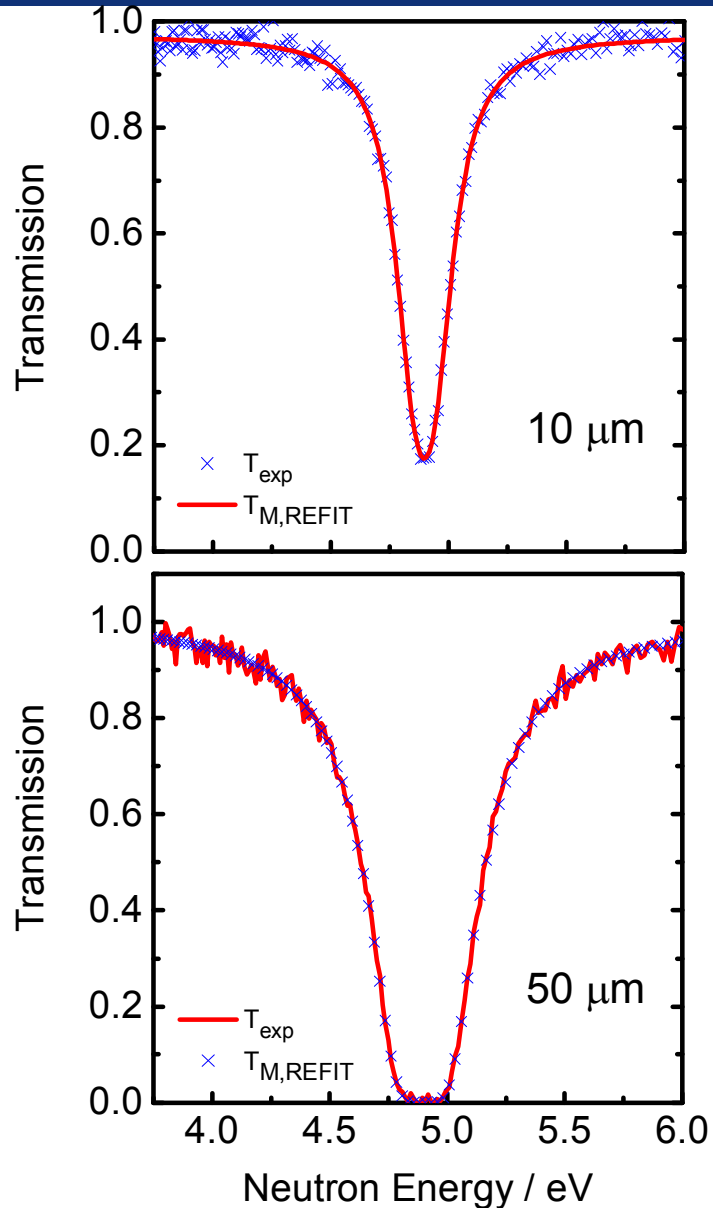
E_r	=	1.260 eV
Γ_n	=	0.464 meV
Γ_γ	=	156.0 meV
g	=	3/4

$\sigma(n,\gamma)$ at 25.3 meV

WF : 142.0 (1.5) b
1.0 % normalization
0.5 % uncorrelated
Counts : 143.3 b

Mughabghab:
145.0 (2.0) b

• **Gold**



$C_{T,\text{exp}}$: only uncorrelated uncertainties due to counting statistics

\Rightarrow

$$\Gamma_n = (15.06 \pm 0.08) \text{ meV}$$

$$\Gamma_\gamma = (121.7 \pm 1.3) \text{ meV}$$

$$\rho(\Gamma_n, \Gamma_\gamma) = 0.55$$

\Rightarrow

$$\Gamma_n = (14.66 \pm 0.30) \text{ meV}$$

$$\Gamma_\gamma = (124.8 \pm 3.7) \text{ meV}$$

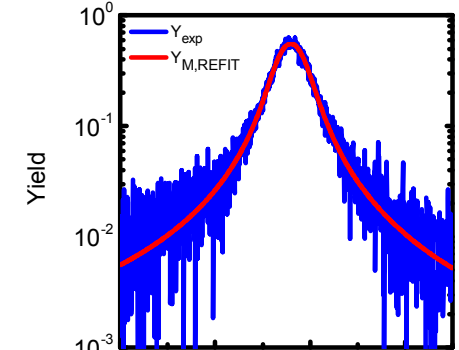
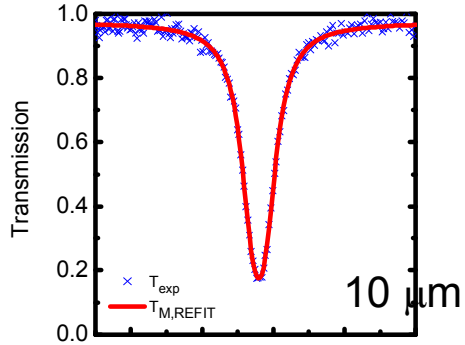
$$\rho(\Gamma_n, \Gamma_\gamma) = -0.96$$

RSA by REFIT

4.9 eV resonance for $^{197}\text{Au}+n$

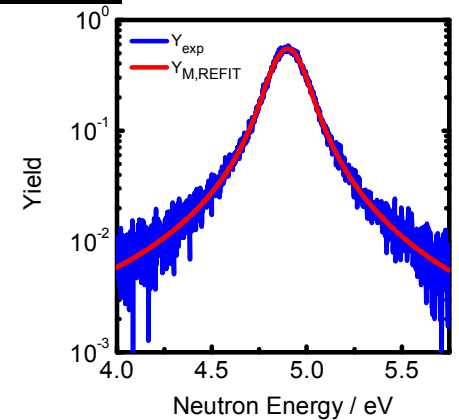
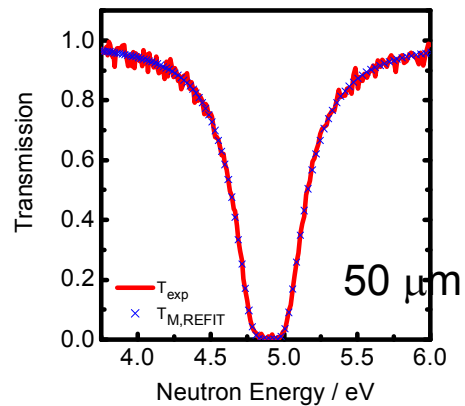
EFNUDAT, Budapest 23/09/2009, P. Sieglar

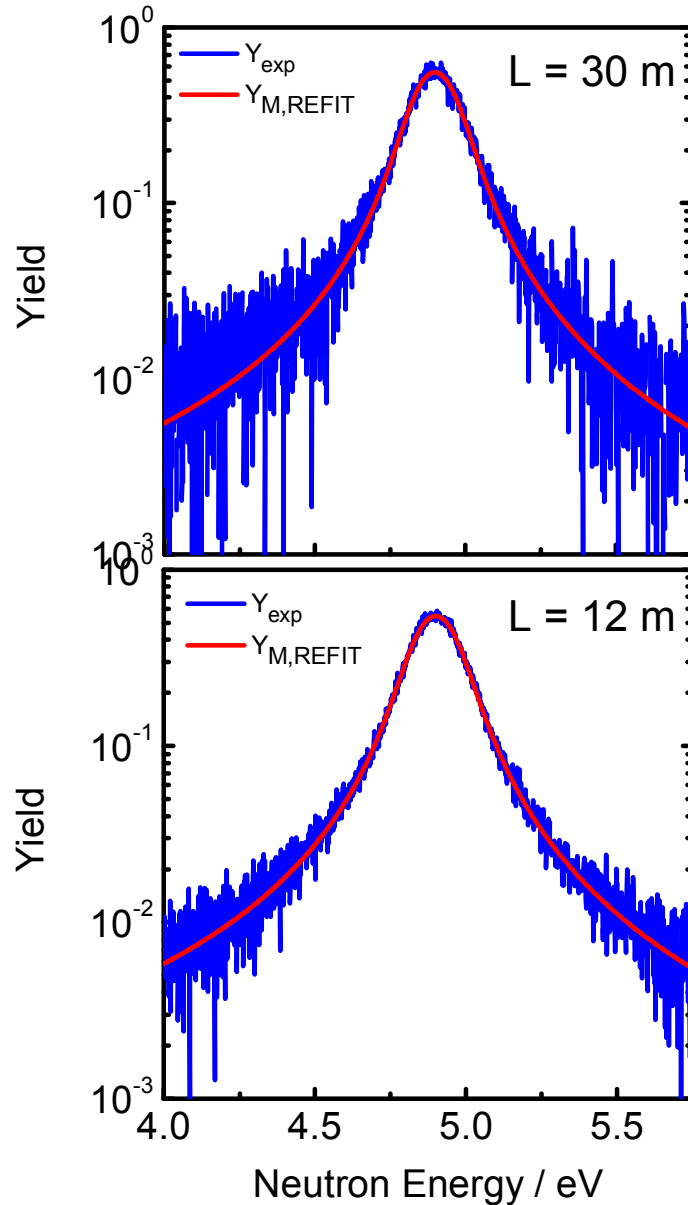
61



Id-number	Measurement	Distance	Angle	Target thickness
(flight path – moderator)				

T1	Transmission	50 m	9°	10 μm
T2	Transmission	50 m	9°	50 μm
C1	Capture	30 m	0°	5 μm
C2	Capture	12 m	18°	5 μm





$C_{Y,\text{exp}}$: only uncorrelated uncertainties due to counting statistics

\Rightarrow

$$\Gamma_n = (15.31 \pm 0.12) \text{ meV}$$

$$\Gamma_\gamma = (118.0 \pm 1.4) \text{ meV}$$

$$\rho(\Gamma_n, \Gamma_\gamma) = -0.50$$

\Rightarrow

$$\Gamma_n = (15.26 \pm 0.15) \text{ meV}$$

$$\Gamma_\gamma = (118.9 \pm 1.2) \text{ meV}$$

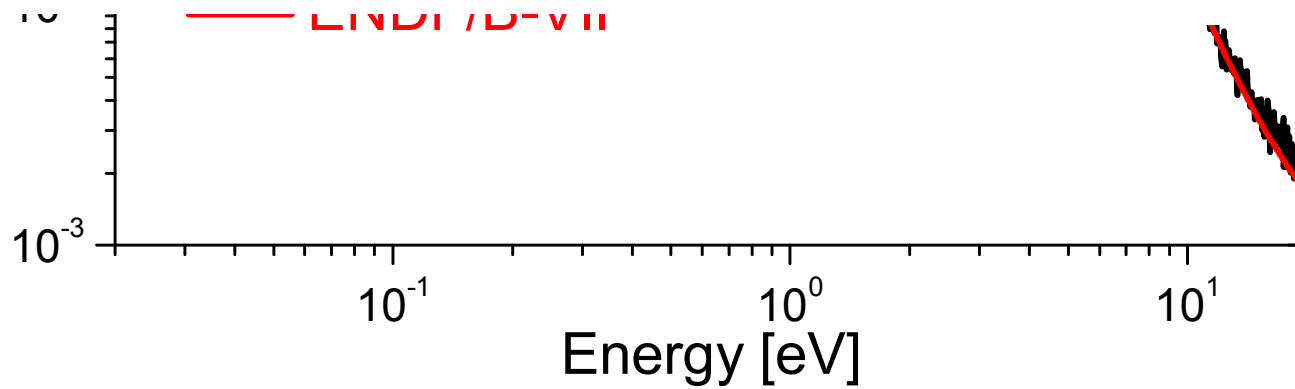
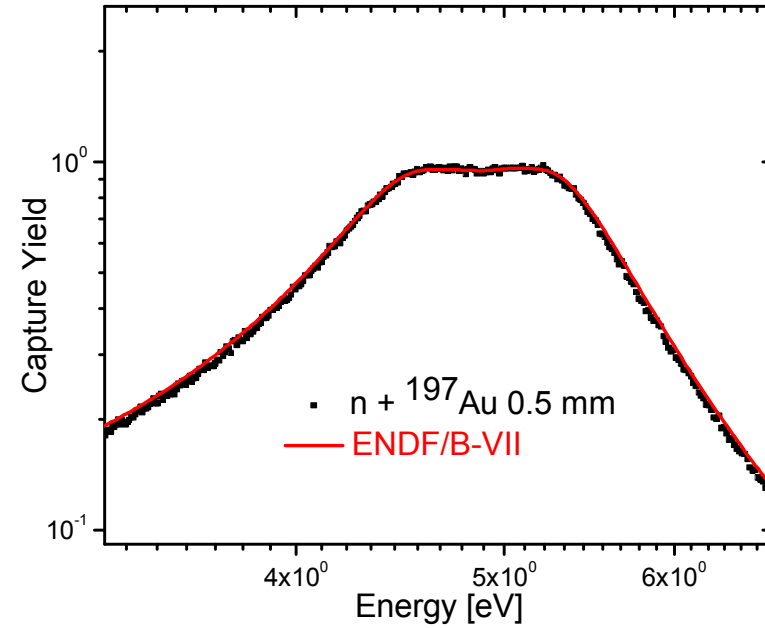
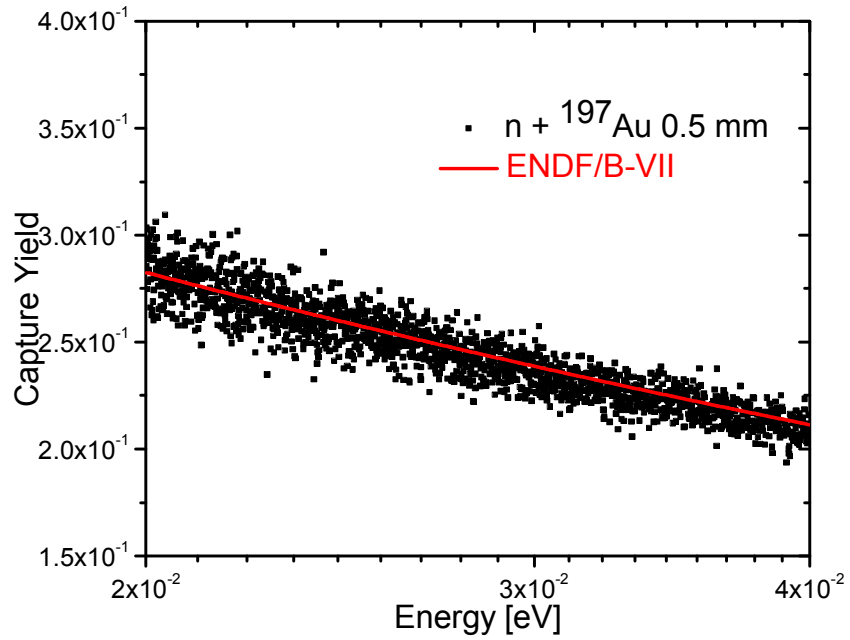
$$\rho(\Gamma_n, \Gamma_\gamma) = -0.63$$

RSA by REFIT

4.9 eV resonance for $^{197}\text{Au}+n$

Measurements	Γ_n / meV	$\Gamma_\gamma / \text{meV}$	$\rho(\Gamma_n, \Gamma_\gamma)$
T1	15.06 ± 0.08	121.7 ± 1.3	0.55
T2	14.66 ± 0.30	124.8 ± 3.7	- 0.96
C1	15.31 ± 0.12	118.0 ± 1.4	- 0.50
C2	15.26 ± 0.15	118.9 ± 1.2	- 0.63
T1 + C1	15.14 ± 0.07	120.0 ± 1.0	0.06
T1 + C2	15.10 ± 0.07	120.2 ± 0.9	- 0.29
T1 + T2 + C1 + C2	15.14 ± 0.06	119.8 ± 0.7	- 0.47

Id-number	Measurement	Distance	Angle	Target thickness
			(flight path – moderator)	
T1	Transmission	50 m	9°	10 μm
T2	Transmission	50 m	9°	50 μm
C1	Capture	30 m	0°	5 μm
C2	Capture	12 m	18°	5 μm



- **14:30-15:00** **Cristian Massimi et al.: The neutron resonance parameters of ^{197}Au from transmission, capture and self-indication measurements at GELINA.**

- **Design and control of experimental parameters is of paramount interest for precise cross section measurements.**
- **Combination of transmission and capture measurements**
- **Dedicated measurements in thermal flux are useful**
- **Cd: capture and transmission final up to 1.5 keV**
- **Rh: measurements finished; analysis: „fine-tuning“**
- **Au: completed; see C. Massimi et al.**